

NEAR-TERM OPPORTUNITIES FOR BEAM-DRIVEN WDM SCIENCE

Richard M. More, LBNL

- **Remarkable phenomena in WDM** <-- *What's there in WDM?*
- **Tools and applications** <-- *Sketch of today's research*
- **Basic questions** <-- *There are many questions*

WDM conditions are easily produced in the laboratory

$$0.1 \text{ eV} < T < 10 \text{ eV}$$

$$0.1 \rho_0 < \rho_0 < 10 \rho_0$$

Surprising that many basic science questions are open

WDM Research

Warm + dense ==> high pressure ==> dynamic experiments

"*Warm*" --> Melt and boil anything, even tungsten

2-phase (liquid-vapor)

Electronic excitation and maybe ionization

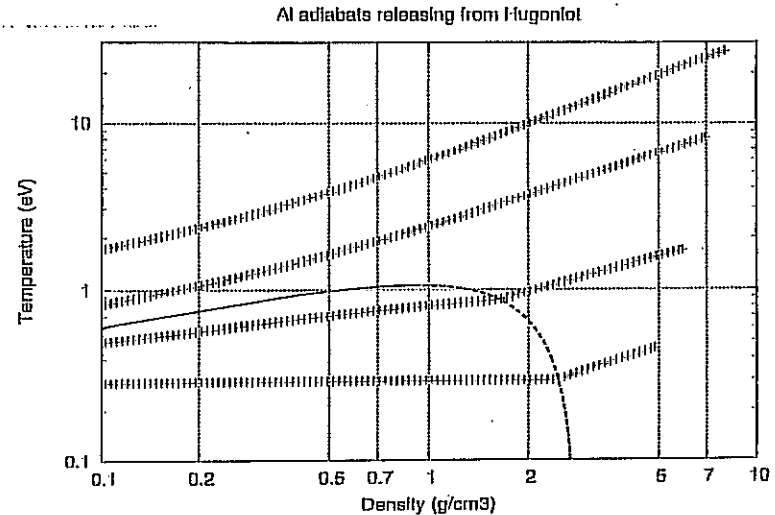
"*Dense*" --> Atoms interact with neighbors (disordered fluid environment)

"*Matter*" --> Material phases and properties

"*Research*" means we want to *learn something new*.

Typical WDM experiments

- A Heat slowly
- B Heat rapidly
- C Shock heat & compress



Time before disassembly

$$\tau \sim L / c_s \sim (1 \mu) / (10^5 \text{ cm/sec}) \sim 10^{-9} \text{ sec}$$

Specific heat

$$C_v \sim 3 \text{ eV/atom-eV} = (2.9/A) 10^5 \text{ Joules/gram-eV}$$

Energy required

$$\text{For } \rho \sim 10 \text{ g/cm}^3 \text{ and } T \sim 1 \text{ eV} \quad E = \rho L C_v T \sim 15 \text{ J/cm}^2$$

Power

$$E / \tau \sim 1.5 \cdot 10^{10} \text{ W/cm}^2$$

For shock launched by a flier:

$$\text{Assume a 5 mm flier, } \rho = 10 \text{ g/cm}^3, v = 10 \text{ km/sec}$$

Energy/area

$$E = 1/2 \rho L v^2 = 2.5 \cdot 10^5 \text{ Joules/cm}^2 \text{ is delivered in } 1/2 \mu\text{sec}$$

Power

$$5 \cdot 10^{11} \text{ W/cm}^2$$

Ticket price

Warm Dense Matter is a distinct subject of study

Existence theorem

USP Laser absorption experiment D. Price et al., PRL 75, 252 (1995)

120 fsec pulse - normal incidence - self-absorption $I = 10^{12}$ to 10^{18} W/cm²

Aluminum is well-described by "usual" plasma theory (QEOS + Drude)

Other materials agree with Al at high-T (plasma) conditions

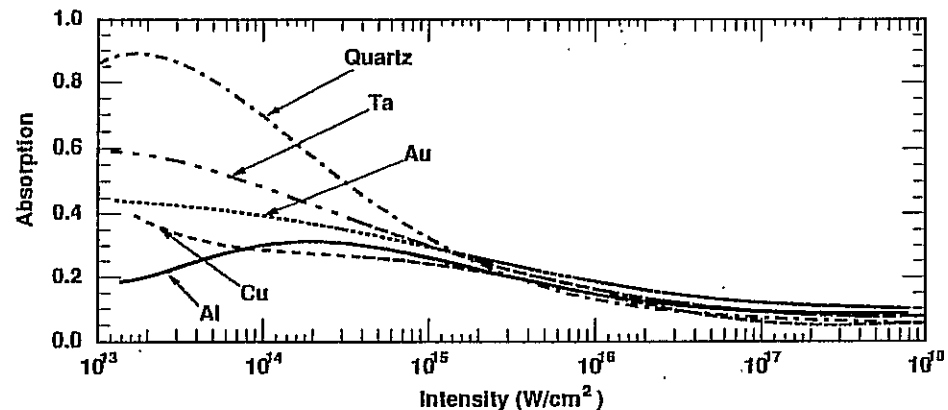
Intermediate range:

Not = solid

Not = plasma

Not a simple interpolation

each material is unique



The experiment shows there are new things to learn.

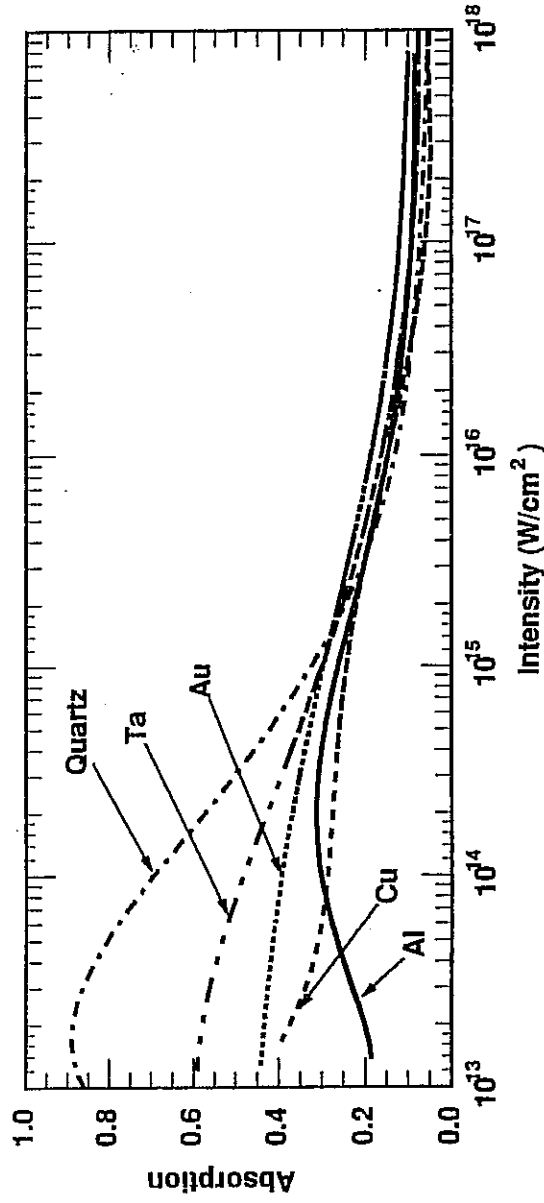


FIG. 1. Absorption fraction vs peak laser intensity for aluminum, copper, gold, tantalum, and quartz targets. In Figs. 1, 3, 4, and 5 laser intensity is the temporal and spatial peak value of the laser intensity.

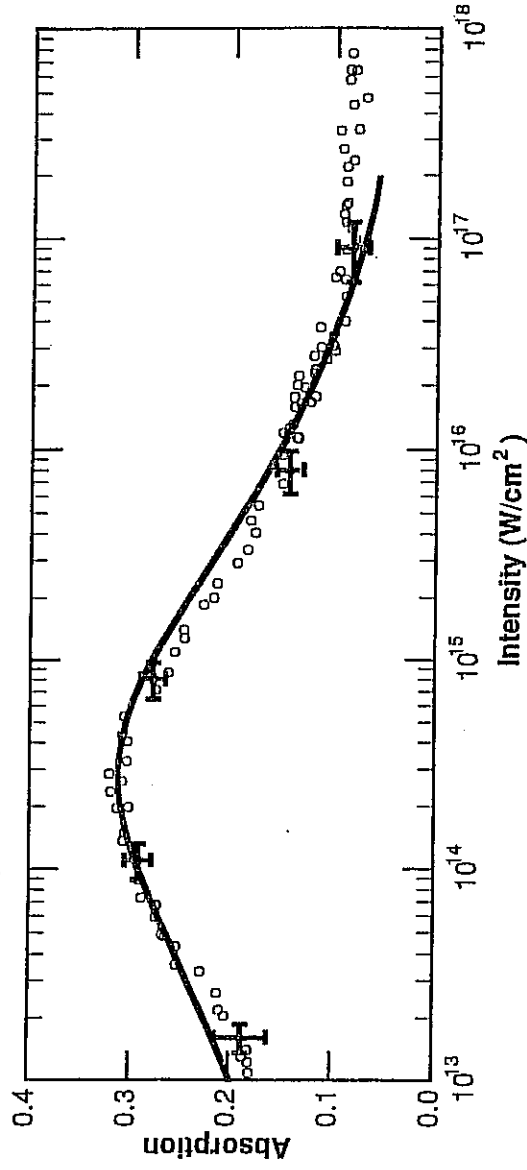


FIG. 5. Measured and calculated absorption fractions for an aluminum target vs peak laser intensity. The calculation assumes a Gaussian laser spatial profile. Error bars shown indicate systematic uncertainty in intensity and random error in absorption. The absolute absorption scale is believed to be known to ± 0.035 .

D. F. Price, R. M. More, R. S. Walling, G. Guethlein, R. L. Shepherd, R. E. Stewart and W. E. White,

Physical Review Letters **75**, 252 (1995).

"Absorption of ultra-short laser pulses by solid targets heated rapidly to temperatures 1 - 1000 eV"

Scientific value of experiments increases with their precision, so good diagnostics are essential.

E_{dep} , expansion $L(t)$, electrical $I(t)$, $V(t)$, optical emission I_ν , XRD

At least four possible methods to measure temperature:

Hydrodynamic release $L \sim 3 c_s(\rho, T) t$ $3 \rightarrow$ ideal gas

Electrical conductivity $\sigma(\nu, T) = \text{current} / \mathbf{E} \text{ field}$

Optical emission $I_\nu = \epsilon(\nu, \rho, T) B_\nu(T)$ $\epsilon \leq 1, \epsilon = a$

X-ray diffraction or scattering (Thomson scattering)

Need to cross-calibrate temperature scales.

Fixed points will help for this.

phenomena that occur in a narrow range of temperatures.

Remarkable phenomena can be predicted

Predictions from preliminary theory and experimental clues

- o \pm ion plasma [$n_e \ll N.$]
- o Metal-Insulator transition [high σ -- $>$ low σ]
- o "Black glass" [$\text{Im}(\epsilon) > \text{Re}(\epsilon) > 0$]
- o Neutral phase of ionic solids ?
- o Mixed valence and shell-crossing [e.g., Pm^{9+}]

Sharply-defined observable phenomenon can help calibrate temperature scales

PLUS-MINUS ION PLASMA

Surprising electronegative plasma predicted for WDM temperature range.

plus and minus ions but $n_e \ll N_+$

Charge density 10^{17} cm^{-3} like semiconductor carrier density

Conduction by charge transfer (n-type, p-type)

Semiconductor-metal transition? Sheath layers?

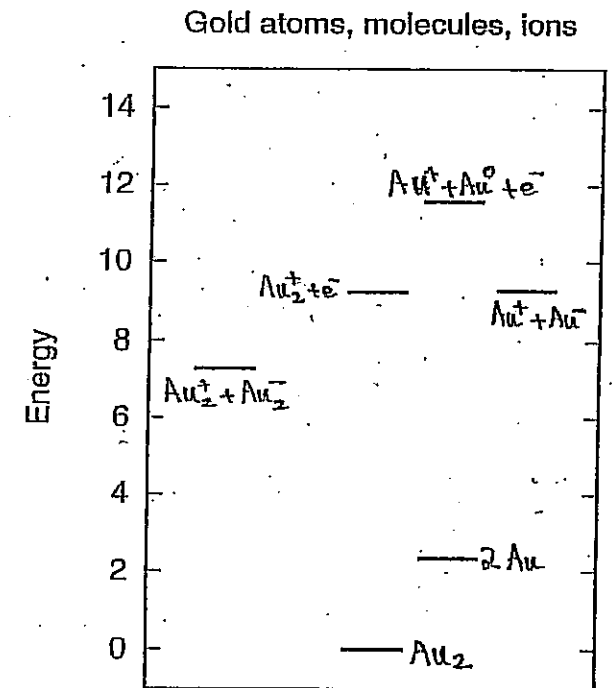
Photodetachment/photoconductivity

Radiation source?

Nonequilibrium heating favors the \pm plasma

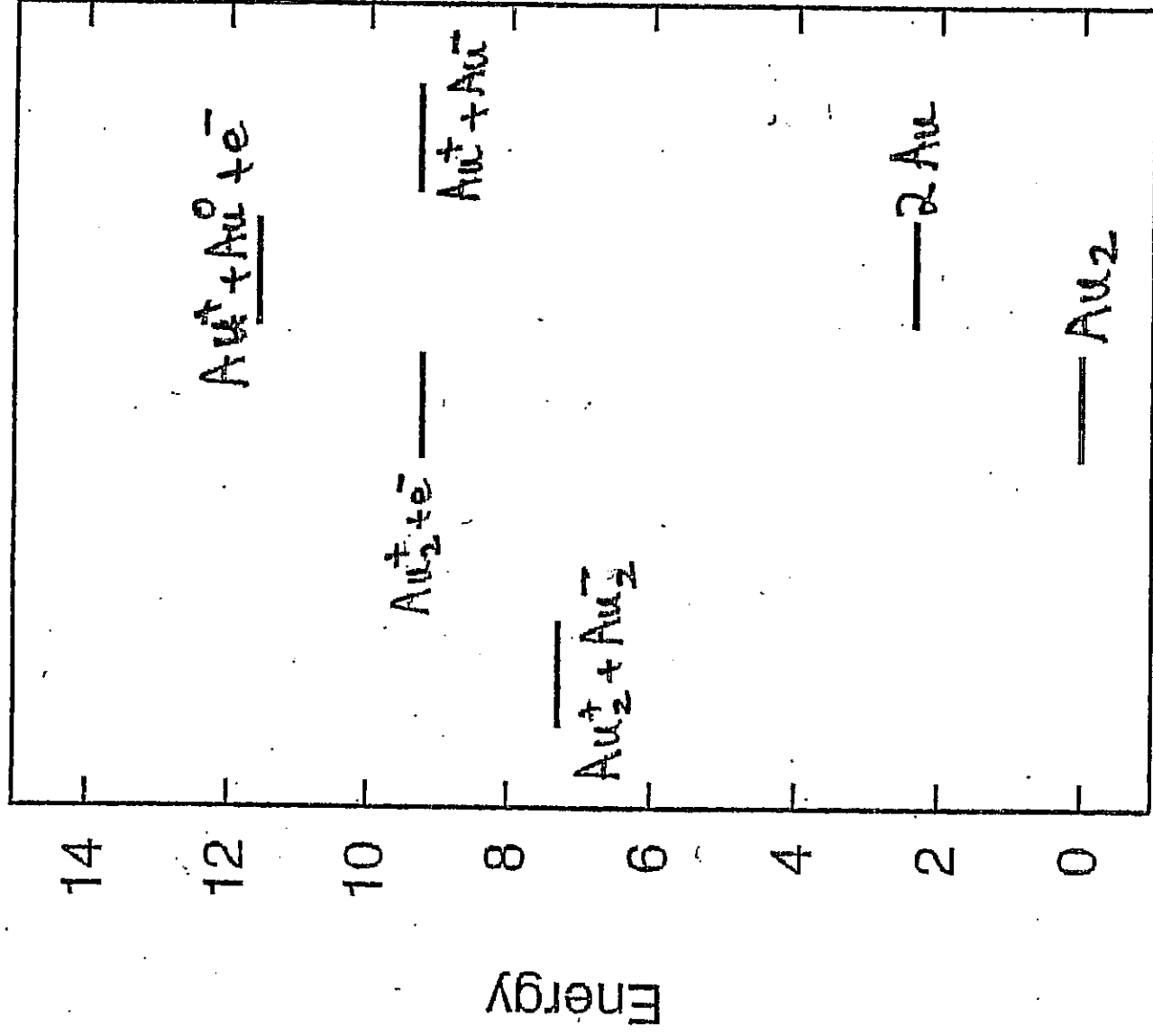
Metals with strong electron affinity (Au, Sn) seem to show this. (Yoneda et al.)

Covalent halogens (Br_2 , I_2) are good candidates (L. Grisham & LBNL group)

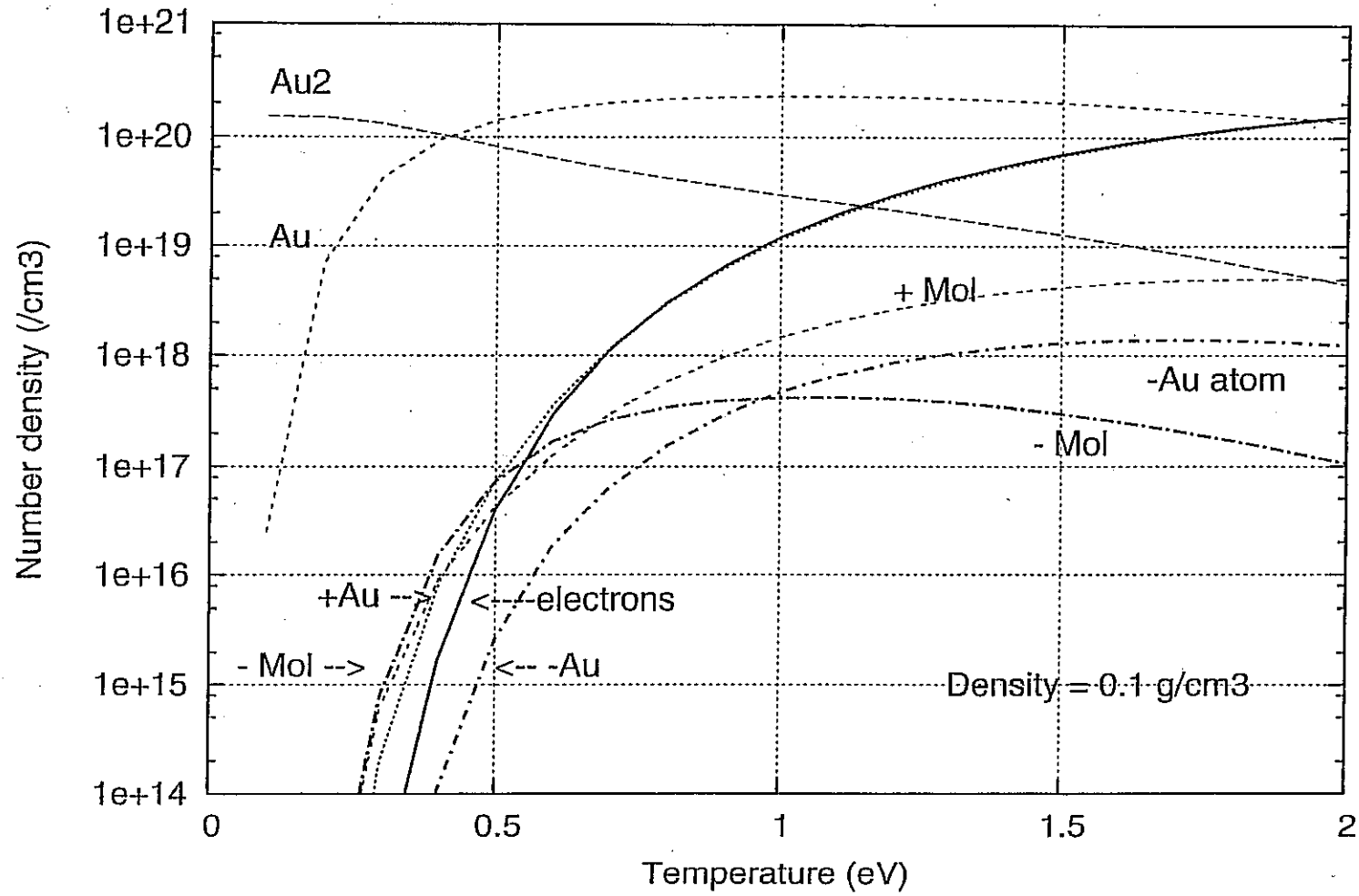


PRL
JOSRT

Gold atoms, molecules, ions



WDM Gold plasma composition



Metal-Insulator transition

[low σ -- > high σ]

SSP: Anderson, Mott, Thouless Datta Libby
Hensel, Faber, Endo, Kitamura
Hg, Cs (low Tc metals) have been studied in detail

WDM: De Silva, Benage, Yoneda - experiments
Laughlin, More, Busquet, Desjarlais (QMD)

At high densities, gap-closing or a percolation threshold.

M. Desjarlais, Cont. Plasma Phys 41 (2001)

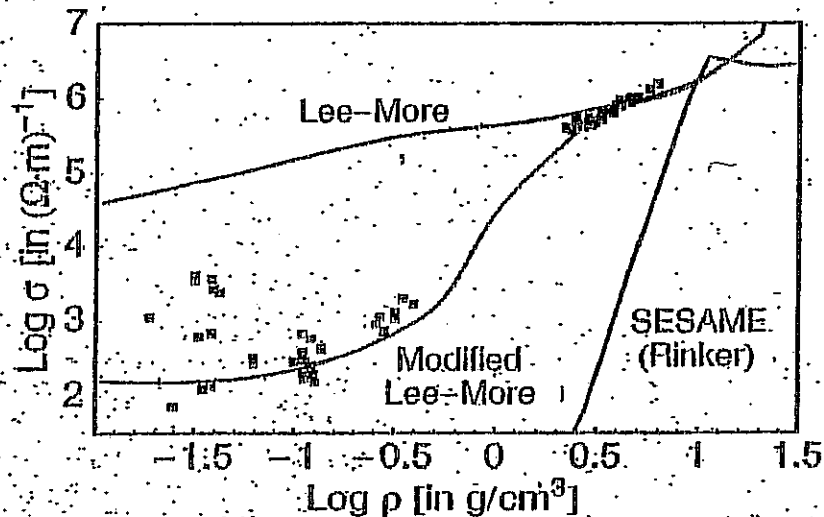
Big questions:

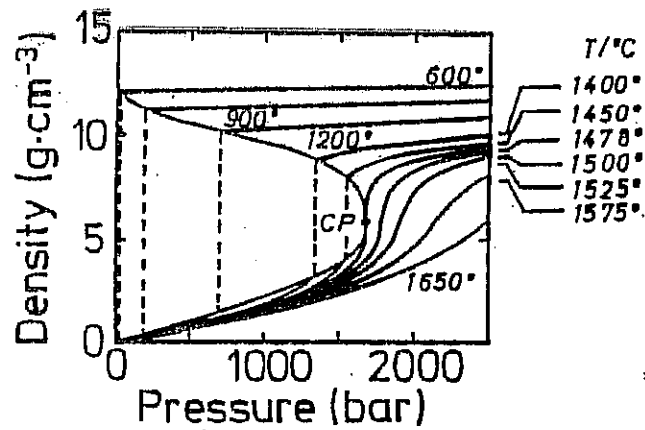
Is there a sharp transition?

Where are the ρ, T boundaries?

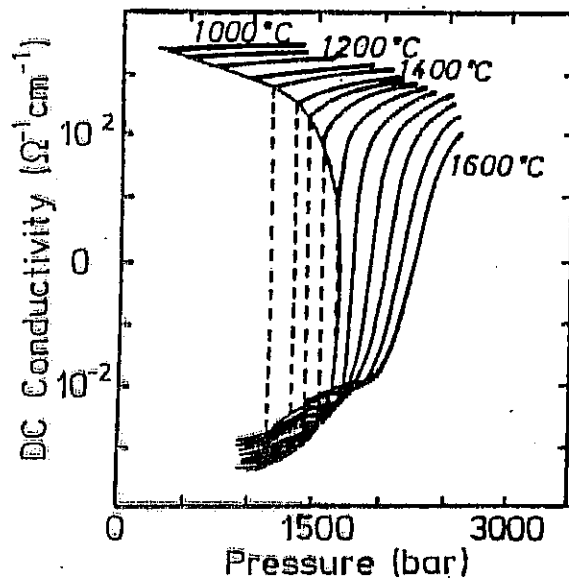
What is the mechanism?

polarons? Mott? Anderson?





(a)



(b)

Fig. 2.4. Isothermal equation of state (a) and DC electrical conductivity (b) data for mercury at sub- and supercritical temperatures as functions of pressure.

2.2 Fluid Semiconductors

2.2.1 Covalent Bonding and Molecular Structure

Fluid metals are not the only fluids in which the thermodynamic state strongly influences the electronic structure. When covalent bonding produces extended network or polymeric structures in the liquid state, the interparticle forces can depend sensitively on the density and temperature. Strictly speaking, the electronic structure of any liquid polymer is state-dependent since the vapor phase consists of monomers or other small species. This state-dependence takes on unusual characteristics, however, if there is significant electronic conductivity. In such a case there is again an electronic transition that accompanies the liquid-vapor transition.

The best elemental example is selenium which melts to form a liquid composed of extended chain molecules (Misawa and Suzuki, 1977). Twofold atomic coordination in these chains is associated with the bonding scheme shown in Fig. 2.5 (see, e.g., Cutler, 1977). The highest occupied electronic states forming a nonbonding or "lone-pair" band are separated from the unoccupied anti-bonding states by an energy gap E_g . It turns out that for selenium, E_g is low enough that a significant number of electrons can be thermally excited to empty states in the upper band or to "acceptor" states associated with "defects" such as broken

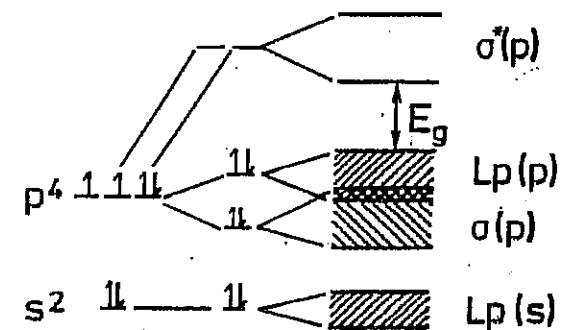


Fig. 2.5. Schematic electronic energy level diagram for selenium showing, from left to right, valence electron states of the free atom, splittings due to interactions with a second atom, and energy bands in the condensed phases of selenium. The p -electron bonding and antibonding states are denoted $\sigma(p)$ and $\sigma^*(p)$, respectively; $Lp(s)$ and $Lp(p)$ denote the lone-pair states for s - and p -electrons, respectively.

"Black glass"

$$\text{Im}(\epsilon) > \text{Re}(\epsilon) > 0$$

SiO_2 Al_2O_3 MgF_2 have Ne-like closed-shell configurations.

Large energy gap (2s-2p to 3s-3p) \implies transparent insulators.

At WDM temperatures, electrons can be excited $2p \rightarrow 3s$

\implies long-lived holes in 2p states

Holes permit intra-atomic absorption $2s \rightarrow 2p$

Experimental indications of this phenomenon:

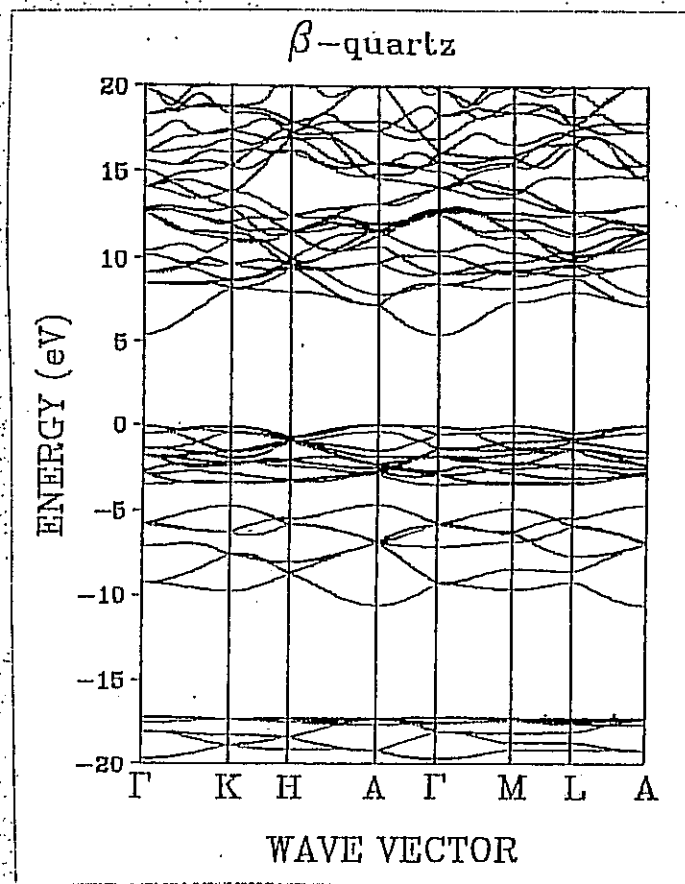
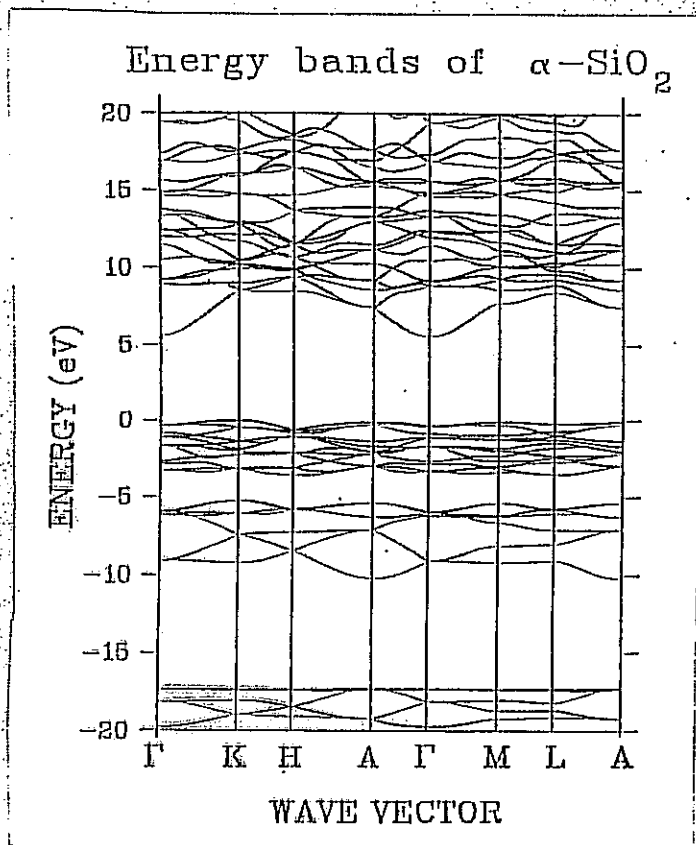
D. Price, Livermore -- USP laser absorption '95

T. Tanabe, Kyushu University -- reactor experiment '03

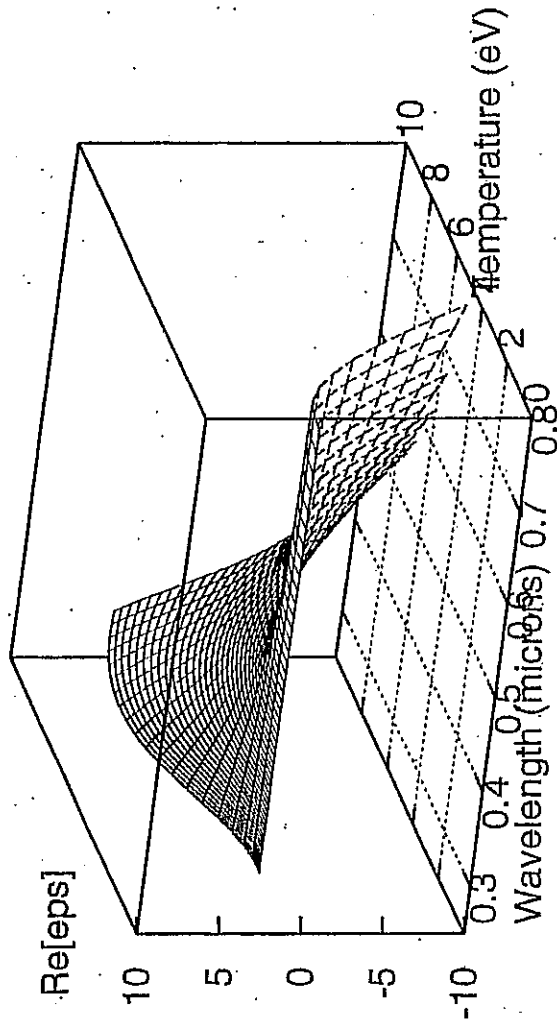
P. Renaudin, CEA Bruyeres -- sapphire window emission '04

P. Lyons, F. Bieniosek - transient darkening of fibers '85

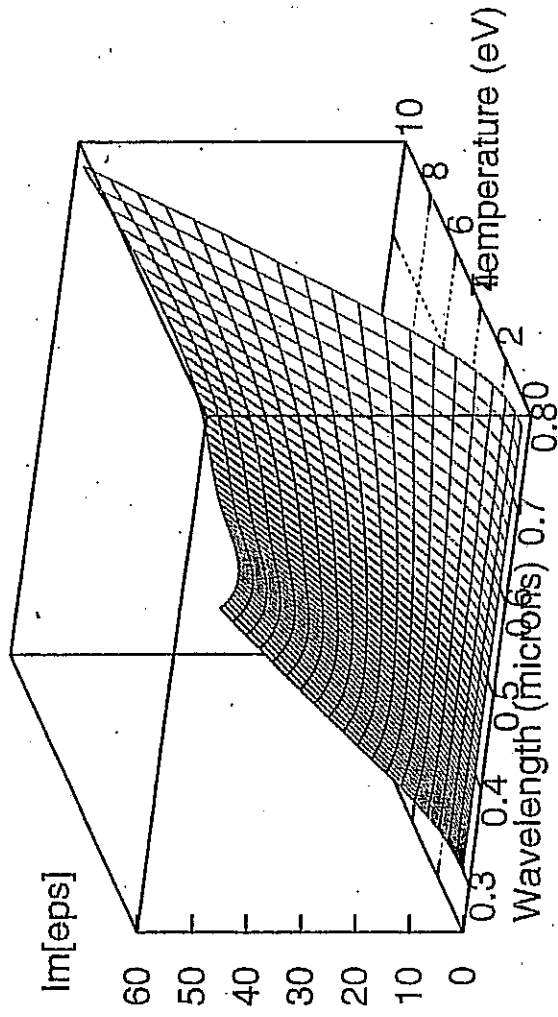
H. Yoneda - Laser and accelerator experiments Dec '06



SiO2 dielectric function



SiO2 dielectric function



Shell-crossing and mixed valence [e.g., Pm⁹⁺]

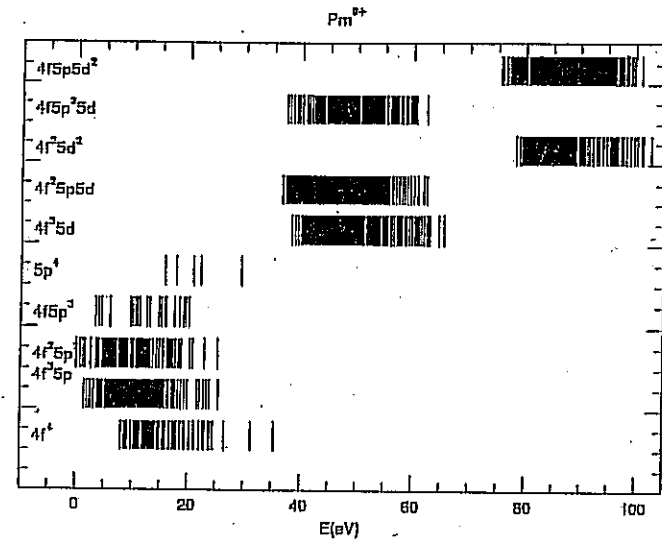
Theory predicts one-electron level crossing in rare earth ions.

Example:

5p-4f crossing in Pm⁹⁺

I. Murakami, NIFS

HULLAC code



Consequences:

Large-scale ("Massive") configuration interaction { 4f^{m-n} 5pⁿ }

Anomalous broadband "white-light" emission

- observed by Dublin laser spectroscopy group

This could set the upper limit of WDM

MODELING TOOLS TO STUDY WDM

R. More, T. Kato, H. Yoneda

J. Wurtele, J. Barnard, G. Penn, A. Friedman

- o **New EOS codes**
- o **Planar hydrodynamic code** "DPC"
- o **Electromagnetic wave-solver** for pump-probe laser ellipsometry
- o **Material models** for Au, Sn, W, glass, Br₂, etc.
- o Monte Carlo simulation of metal-insulator transition
- o **Foil heating simulations** using LLNL Hydra code. Foam targets.
- o **Beam-target interaction** codes for ion beam propagation and heating

Warm dense matter and the liquid-vapor transition

1st generation: Hydrodynamics with equilibrium EOS

Two-phase EOS with Maxwell construction

2nd generation: Evaporation kinetics, surface tension, yield strength

Evaporation / condensation

fast at high ρ , high T

slower at low ρ , low T

"Spinodal" region has a micro-instability when $\frac{\partial p}{\partial \rho} < 0$

For hot metals, EVERYTHING should be tested by experiment.

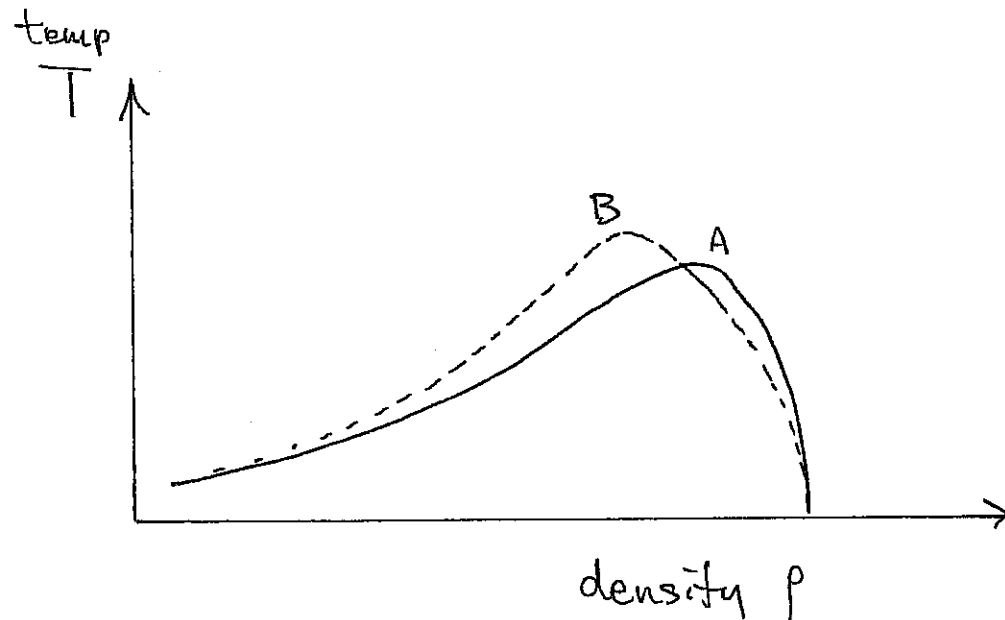
QUESTION:

Do excitation and ionization distort the 2-phase boundary ?

Tungsten is a candidate:

$$T_c = 1.1 \text{ to } 1.5 \text{ eV}$$

$$\text{IP} = 7.98 \text{ eV}$$



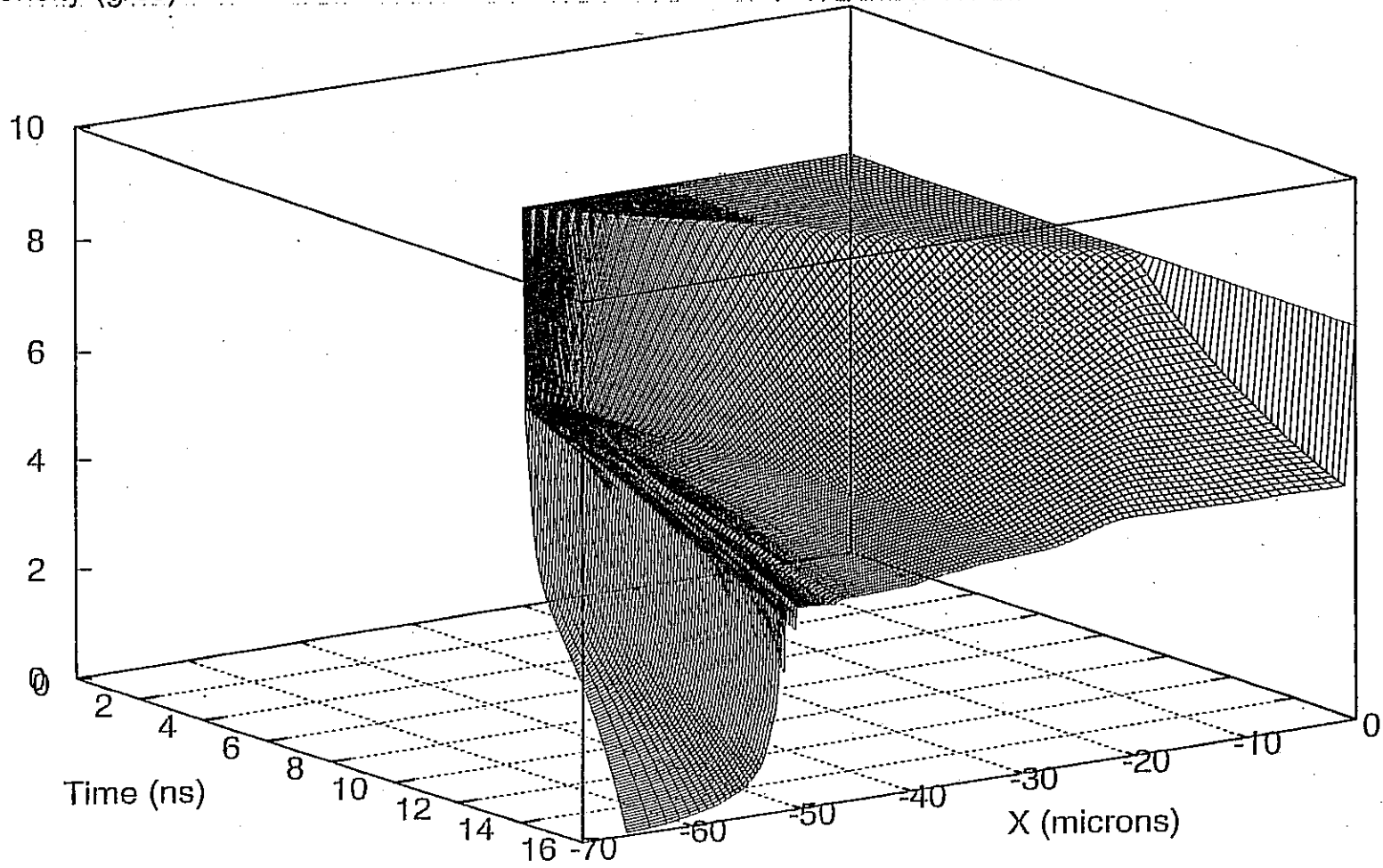
New Tungsten EOS constrained by QMD data from M. P. Desjarlais (SNLA)

Atom pair-potential $V(R)$ changes with excitation, ionization.

For these materials, the fluid "corresponding states" EOS is invalid.

Sn release from solid-density ($T_0 = 1 \text{ eV}$)

Density (g/cc)



Hydrodynamic release into liquid-vapor region:

accurate EOS + good resolution + robust EOS look-up

Release from a uniform initial state: More, Kato, Yoneda find *shelf structure*.

- o Sharp interface of liquid-metal region
- o Low-density ($2-\Phi$) precursor

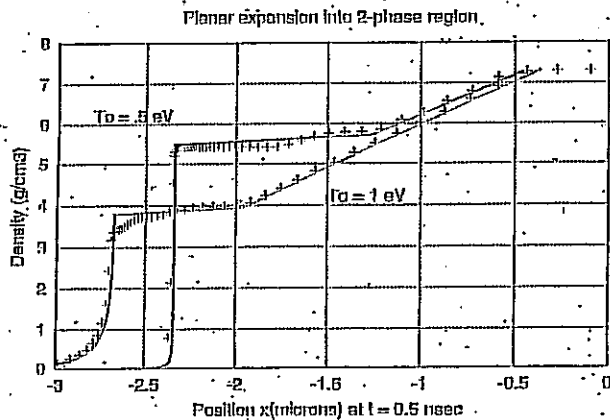
Numerical simulation *agrees with* exact analytic solution

(*) also agrees with Nishihara particle simulation

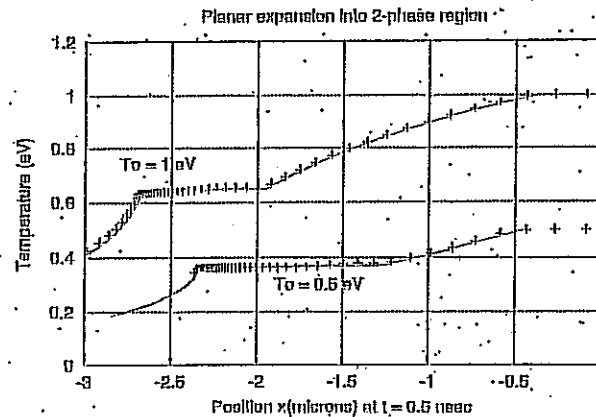
----- = Exact solution

+---+---+ = Numerical hydrodynamics

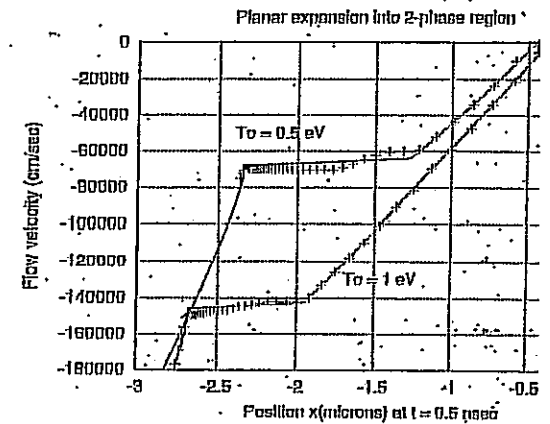
DENSITY



TEMPERATURE



FLOW SPEED



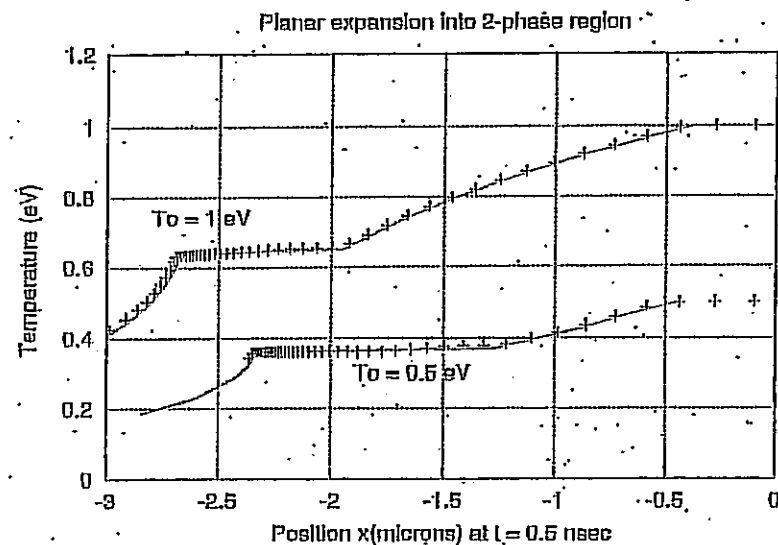
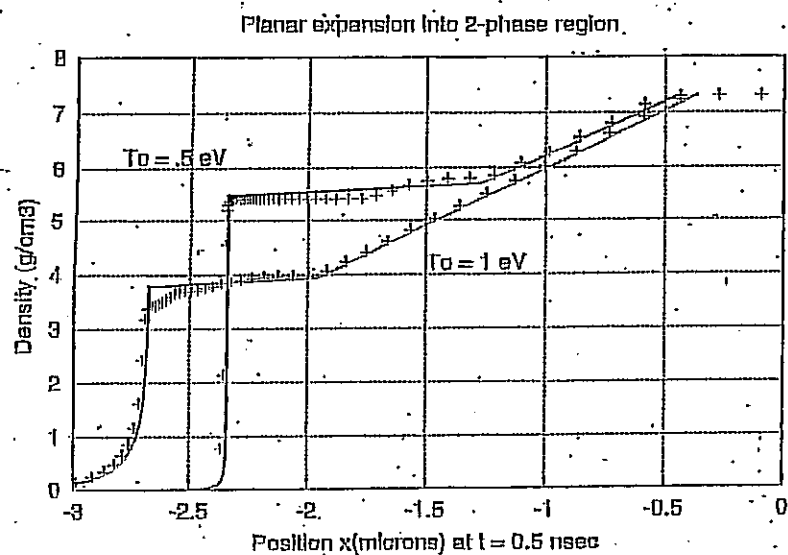
Numerical simulation agrees with exact analytic solution:

----- = Exact solution

+ + + + = Numerical hydrodynamics

DENSITY

TEMPERATURE



Many applications of WDM information:

Fusion research -

Pulsed power - electrical heating and wire dynamics

ICF - early-time behavior; scale-up pathway for HIF; debris

MFE - Pellet injection dynamics; wall interactions

Radiation sources - EUV

Laser cutting and machining

Plasma processing, high-temperature chemistry

Future high-power semiconductors?

Geophysics, astrophysics, condensed-matter

MAKING WARM DENSE MATTER:

Electric current heating

(since 1920)

I(t), V(t) -- > know energy input
Energy in WDM range
Inexpensive sources

rise time > 10 - 50 nsec
expansion -- > nonuniform
conductivity, deposition change

Laser heating

(since 1970)

Wide range of temperatures
Short pulse -- > high density
High rep-rate -- > data rich
Fast diagnostics (pump-probe)
Large and small facilities exist today

Deposition at low density
Absorption changes with T
Needs hydro modeling
nonequilibrium
competition for laser time

Accelerator heating

(today)

Uniform deposition over ~ 1 micron
Robust deposition, any target
Foam targets with uniform deposition

Power limit in today's machines
Need space-time focusing
Need diagnostics

X-ray heating

(future)

Real discoveries will be verified several ways

MATOME

In WDM there are new things to learn

New phenomena

limited range of T == > temperature flags

Technologies for WDM research:

Heating - X-rays, shock release, lasers, electrical, e-beam, and accelerators

Diagnostics - hydro speed, electrical, optical, UV and X-ray transmission

Modeling, theory - Hydro codes, EOS, conductivity, atomic kinetics, QMD

WDM research is "in construction"

Basic phenomena

Basic organization