

Efficient Laser Ion acceleration in an Intense-Short-Pulse-Laser Foil Interaction

- Ion Beam Radio Therapy
=> Compact Ion Accelerator

- 1) Collimated Ion Beam by a Multi-Hole Target
- 2) Possibility of Energy Spectrum Control

US-Japan Workshop September 7-8, 2009, San Francisco



Contributors & Friends

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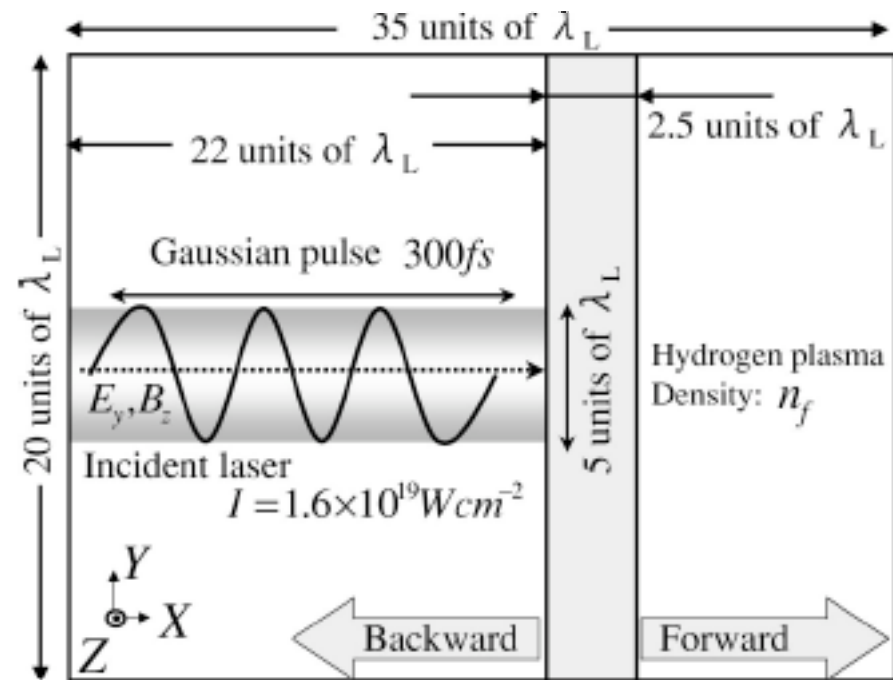
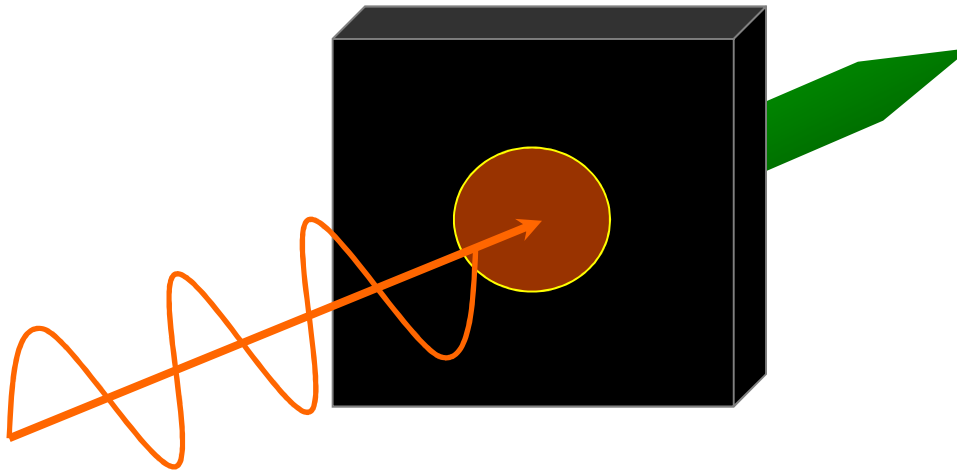
J. Limpouch, O. Klimo

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A. Andreev

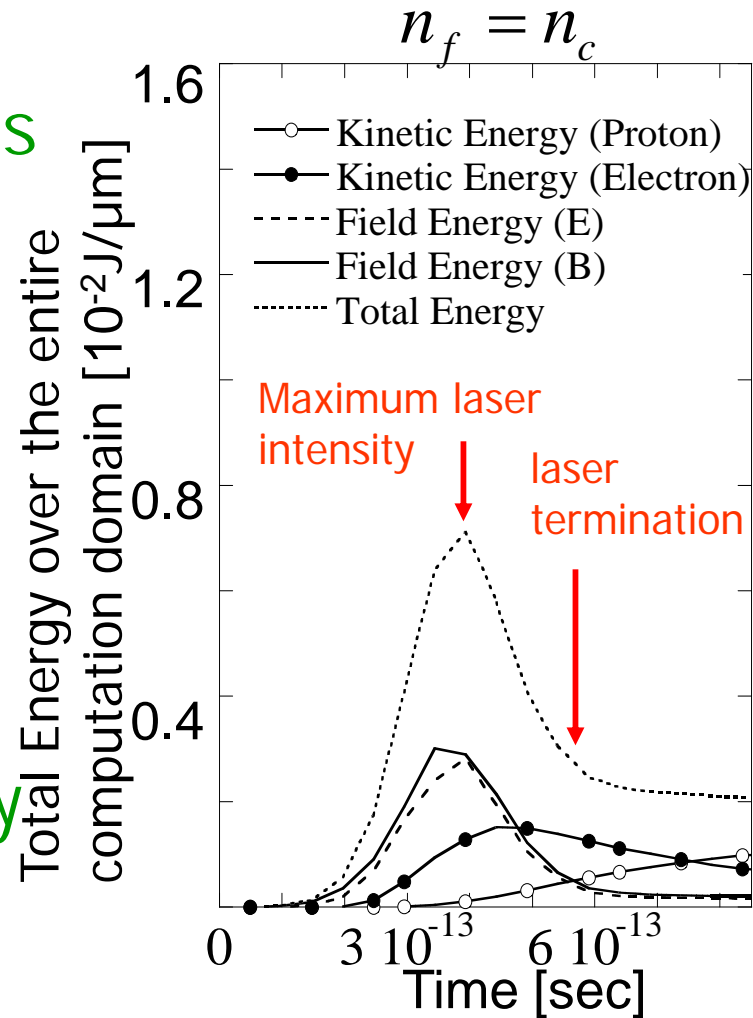
S.I. Vavilov State Optical Inst., St. Petersburg, Russia

Laser Ion Accelerator



Continuous Increasing of the Electron Energy in Optimal Density

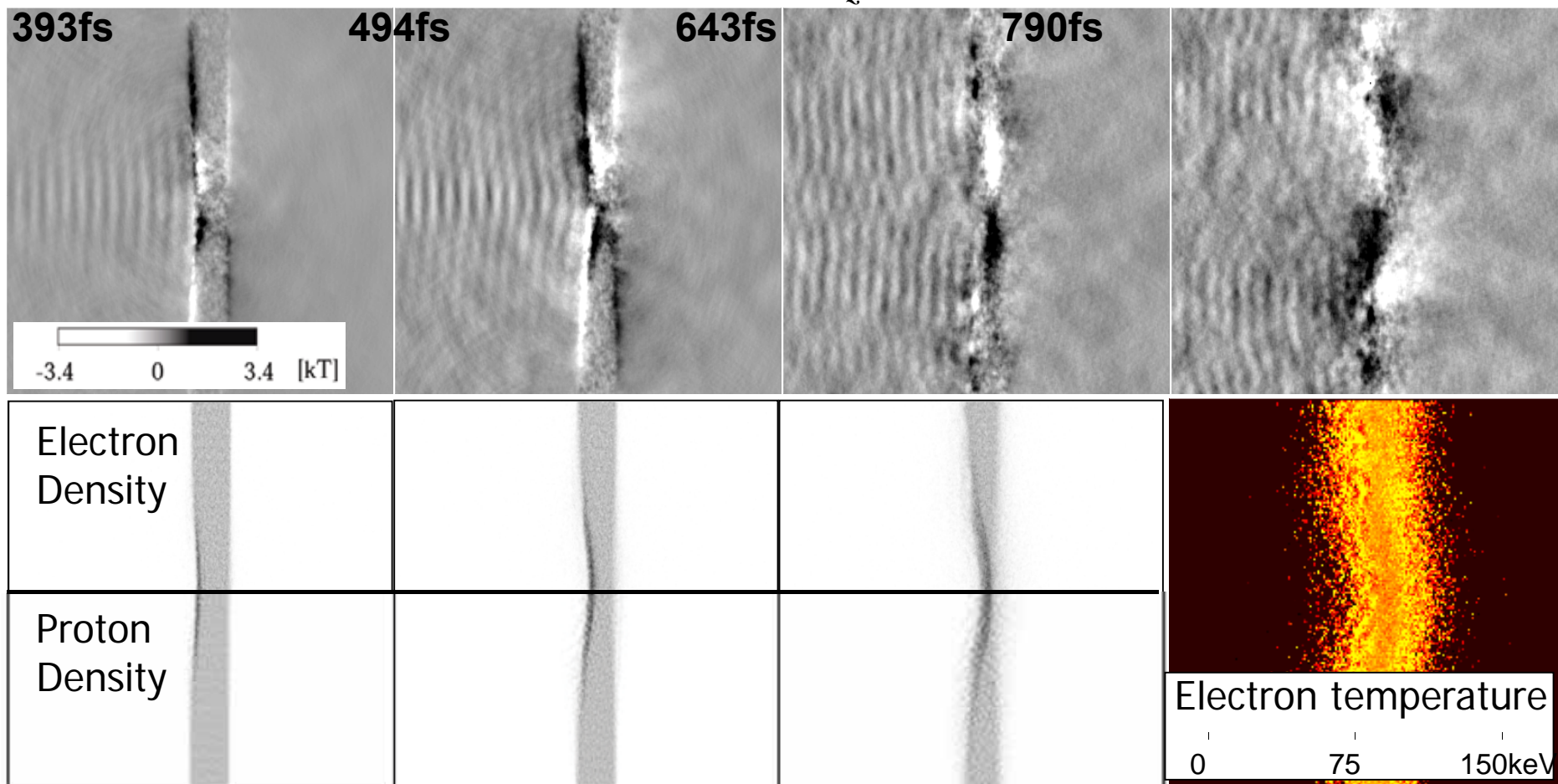
- Energy Storage as Excited Strong Electric and Magnetic Field.
- Continuous Energy Transfer to Electrons from the Stored Field Energy



Excited magnetic field ($n_f=2n_c$)

Averaged magnetic field in one laser period.

$$B_z^{avg}$$



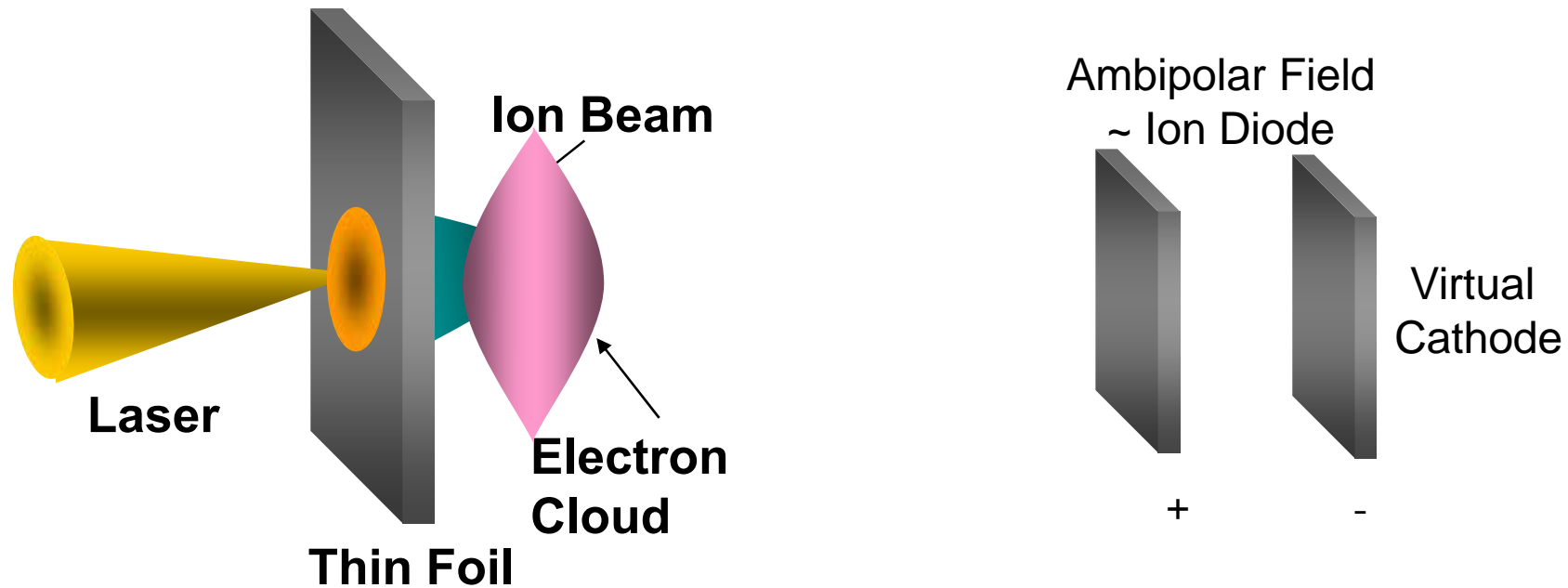
Issues of Laser Ion Accelerator

1. Ion beam quality
 - transverse divergence
 - energy spectrum control
2. Low energy efficiency from laser to ion beam
 - ~ a few % or less
3. Total number of ions accelerated
 - ~ 10^{12} particles or so
4. Low laser efficiency

Problems of Laser Ion Accelerator

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Collimated Ion Beam



Origin of ion beam divergence:

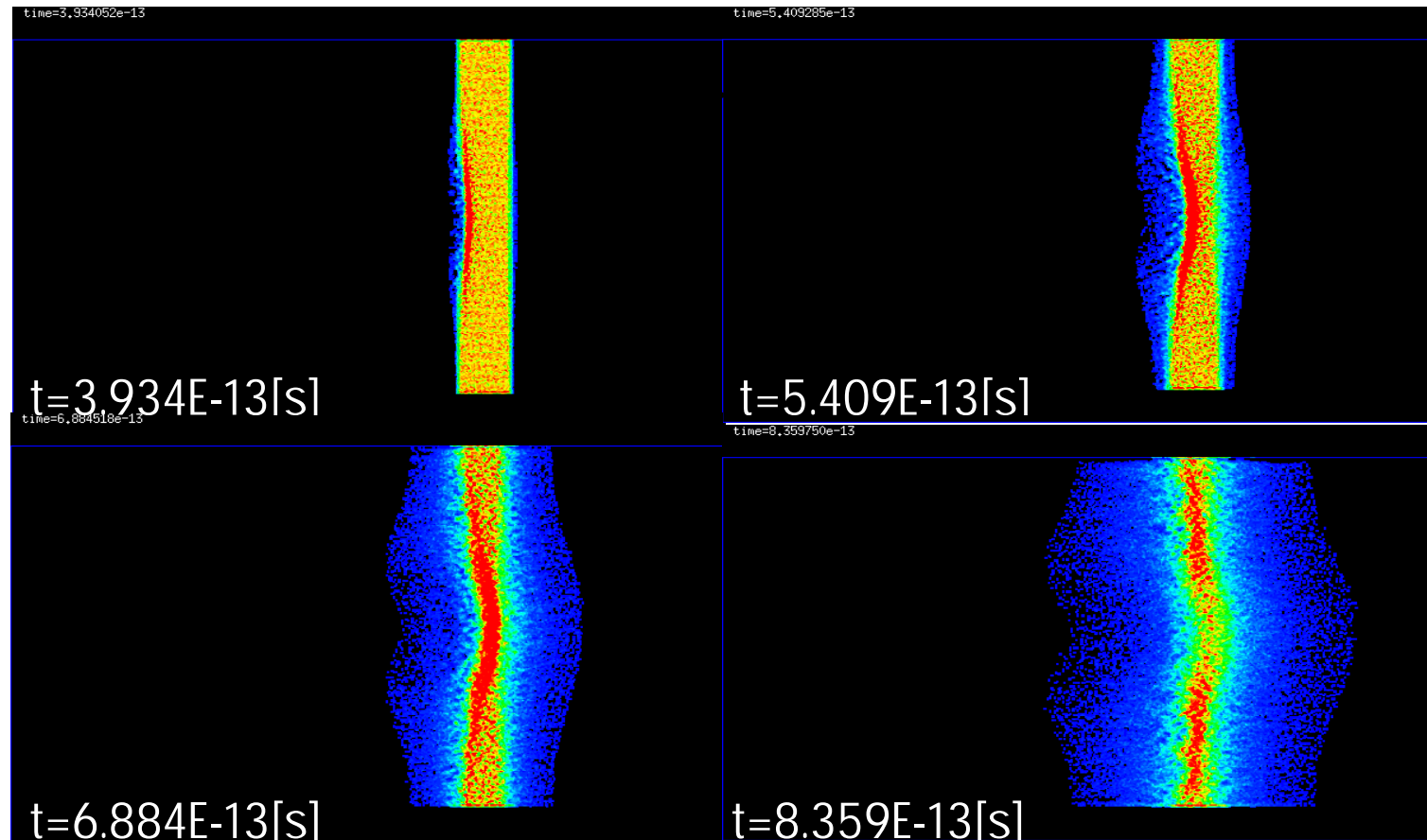
- / Edge fields of ion source & electron clouds
- / Ion beam temperature

Suppression of transverse proton divergence
by shielding edge fields of electron cloud & ion source.

Simulation Results

Target density $n_F = 2 \times n_c = 2.01 \times 10^{21} [cm^{-2}]$

Proton

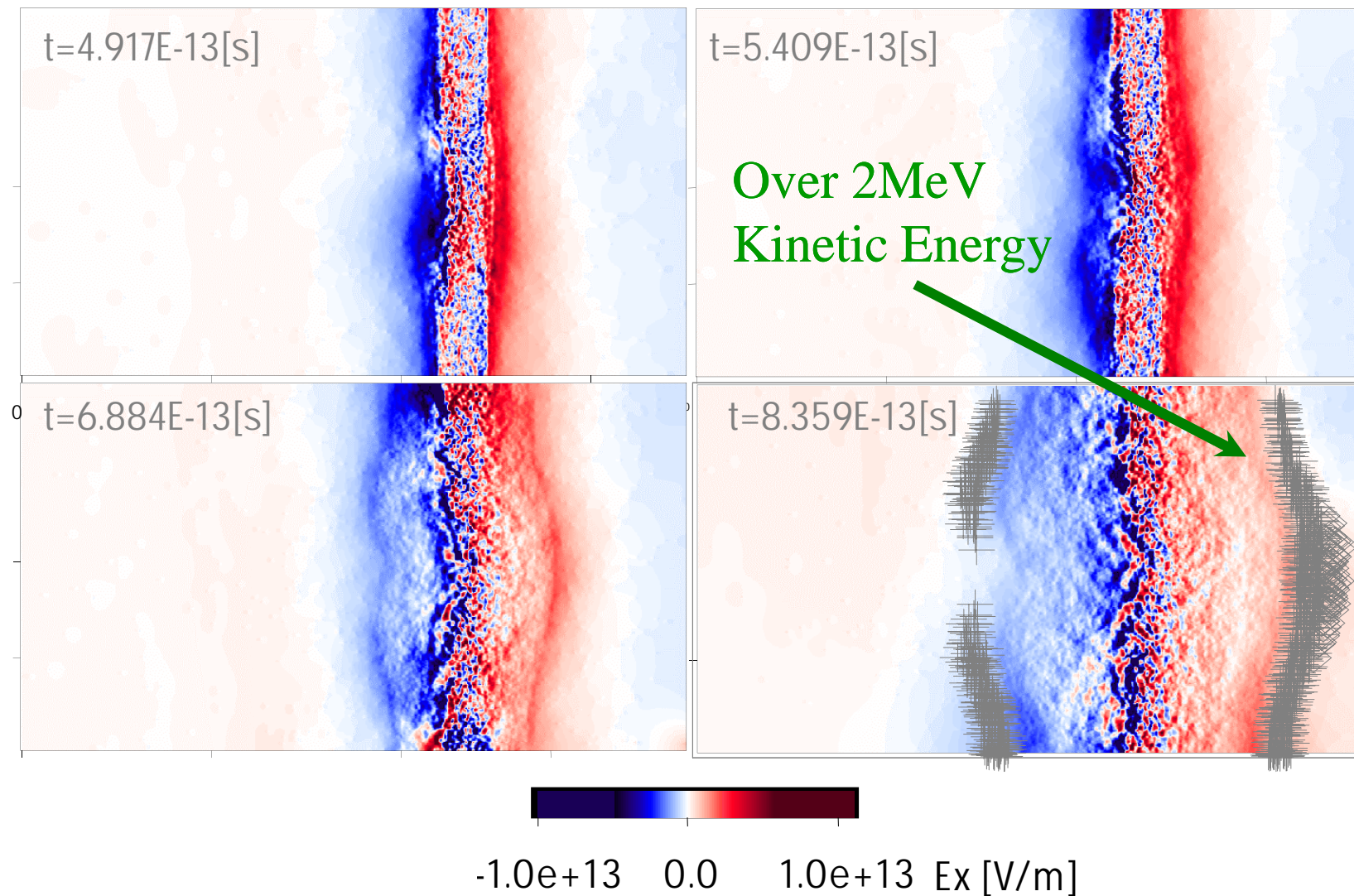


Maximum Kinetic Energy of Proton

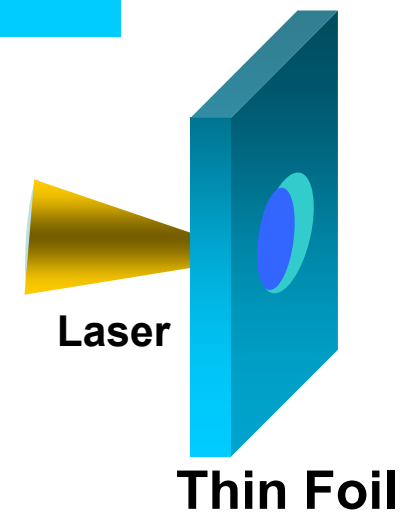
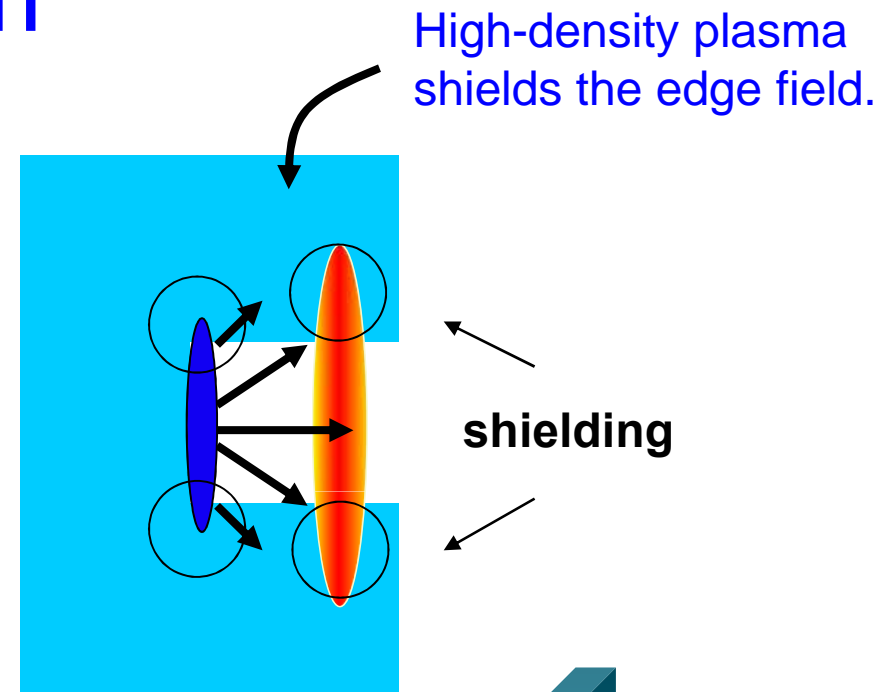
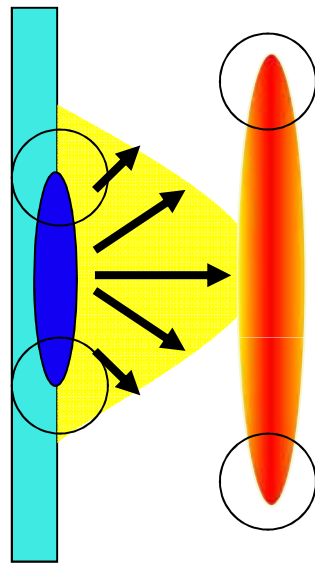
Forward: 3.73[MeV] Backward: 2.89MeV

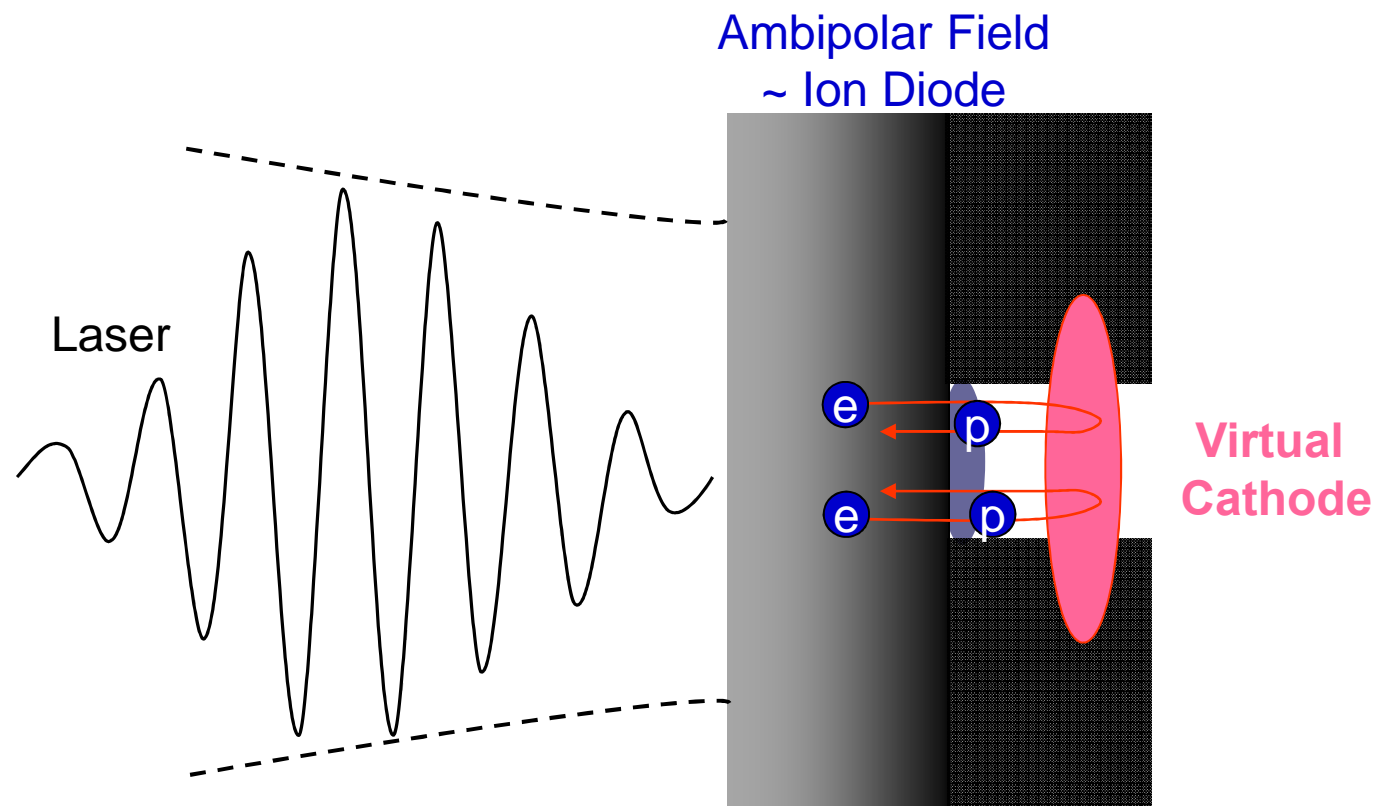
Simulation Results

Electrostatic Field (Time developmet of Ex)

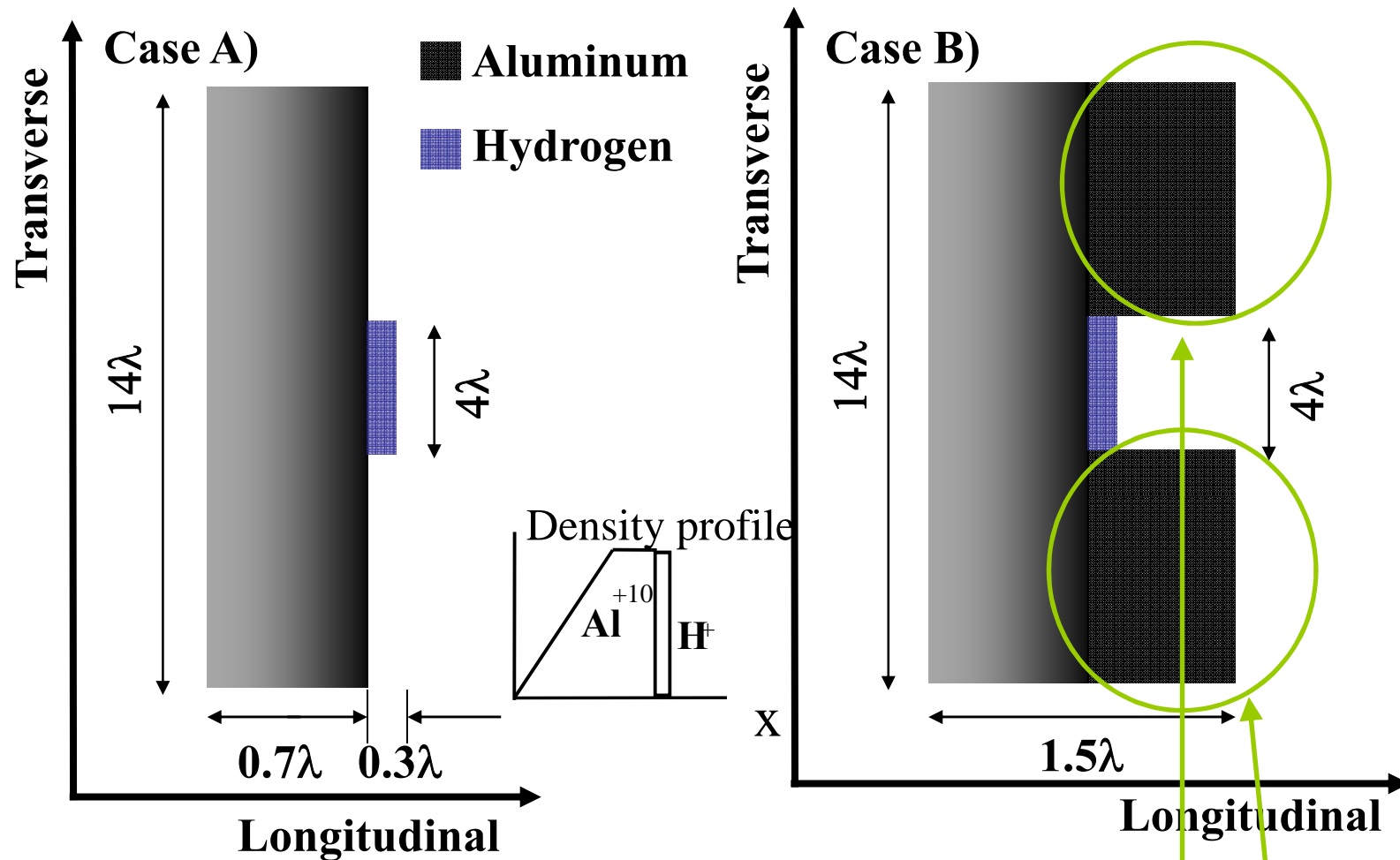


Collimated Ion Beam





Initial Target Profile

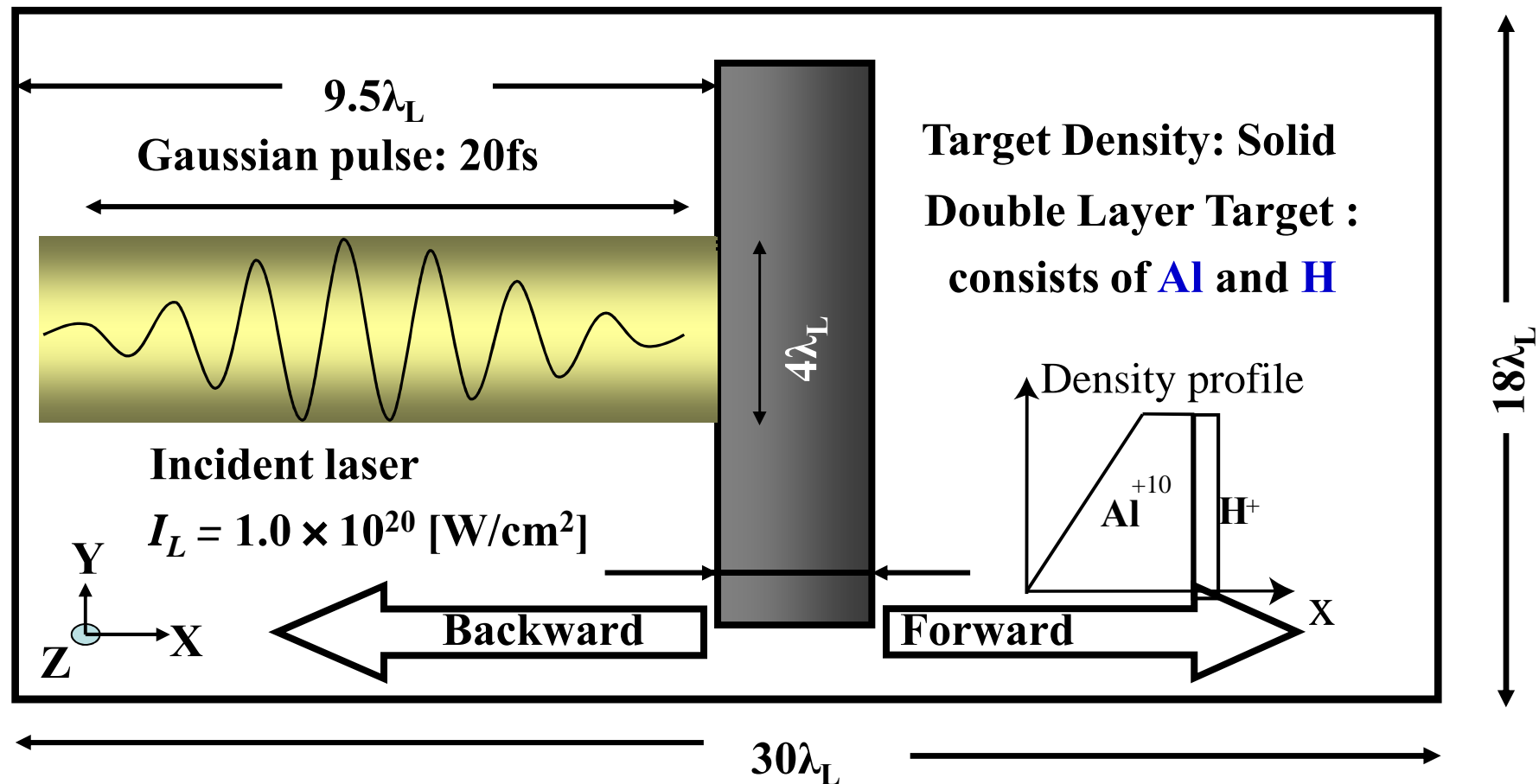


Target of Case A is the conventional slab foil plasma.

Target of Case B has a hole at the opposite side of the laser illumination.

Gradation part in figures mean density gradient.

Simulation Model of 2.5D PIC Simulations



Initial Parameter Values

- **Intense laser pulse**

Wave length: $\lambda_L = 1.053[\mu\text{m}]$

Gaussian laser duration: $t_L = 20[\text{fs}]$

Laser intensity: $I_L = 1.0 \times 10^{20}[\text{W}/\text{cm}^2]$

Laser spot diameter: $r_L = 4\lambda_L$ (FWHM)

- **Target**

Double layer target consists of Al and H

Initial density: solid

Initial distribution: Partial balance-Maxwell

distribution (temperature $T_e = 1.0, T_i = 1.0[\text{KeV}]$)

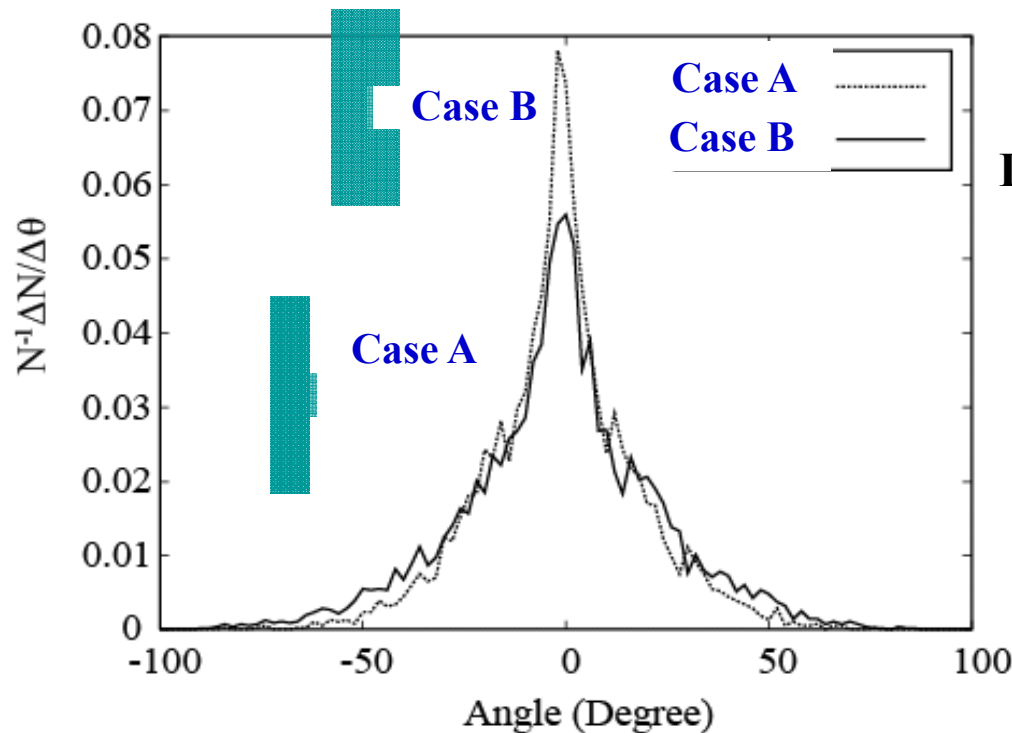
- **Calculation conditions**

The calculation mesh size $\Delta x = \Delta y$ is 0.02λ

The integration time step Δt is $0.04 \times \Delta x/c$

We employ about 1.6-million super particles in our simulations

Estimation of Divergence Angle



Divergence Angle:

$$\frac{\sqrt{\sum_{i=1}^N (\Delta\theta)^2}}{N}$$

Case A : 2.805

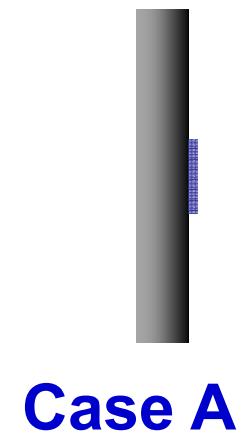
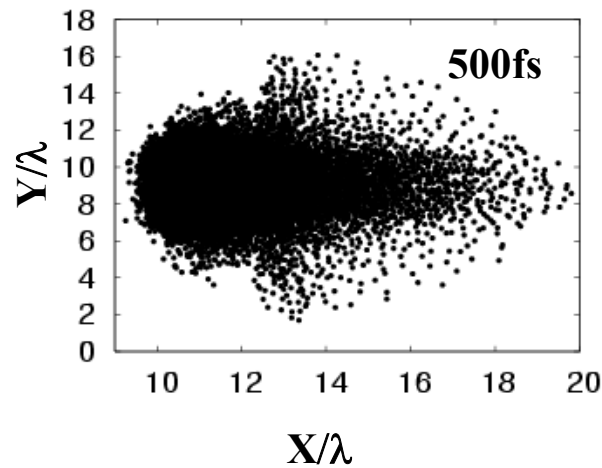
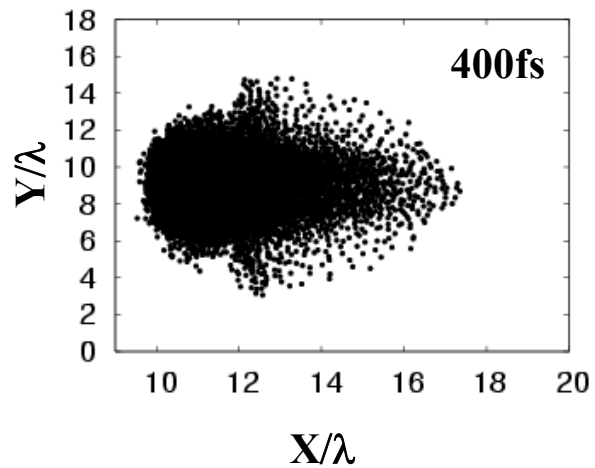
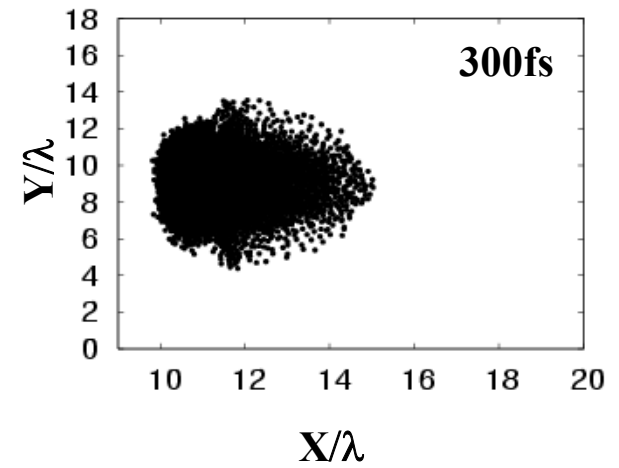
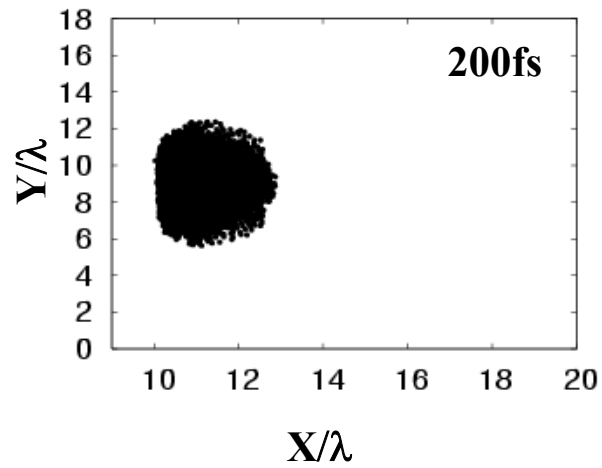
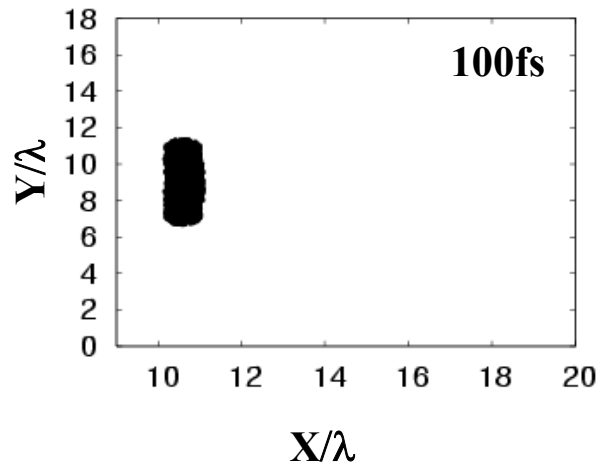
Case B : 2.399

N: total particle number
of the proton beam.

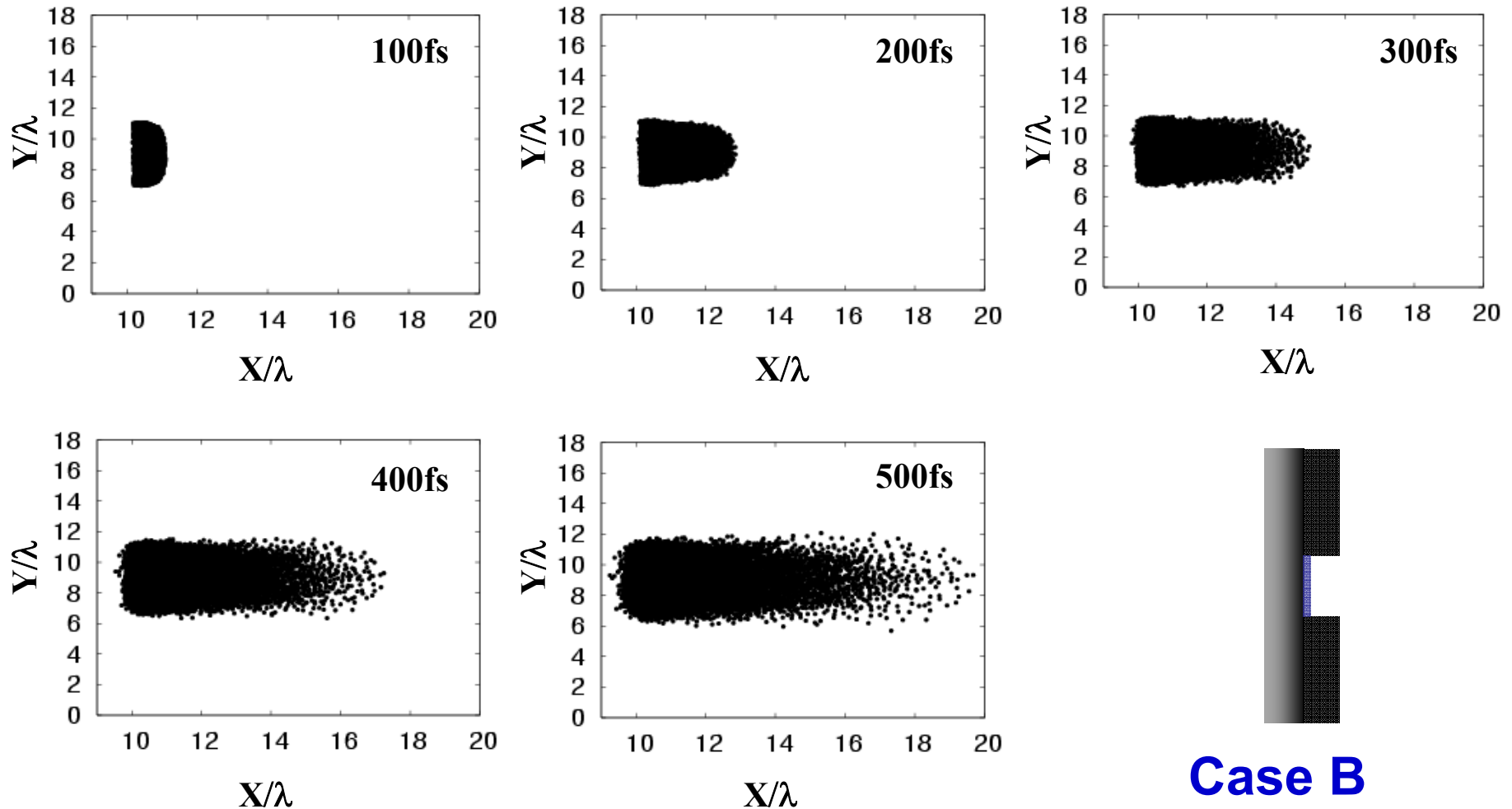
$\Delta\theta = \arctan(v_y/v_x)$.

These values also mean that the proton beam divergence in case of the target with the hole is suppressed compare with that the case of the conventional slab target. The proton beam quality in Case2 is controlled by the electron cloud localization, which is realized by the configuration of the plasma target (Case2).

Proton Distribution in X-Y Space in Case A

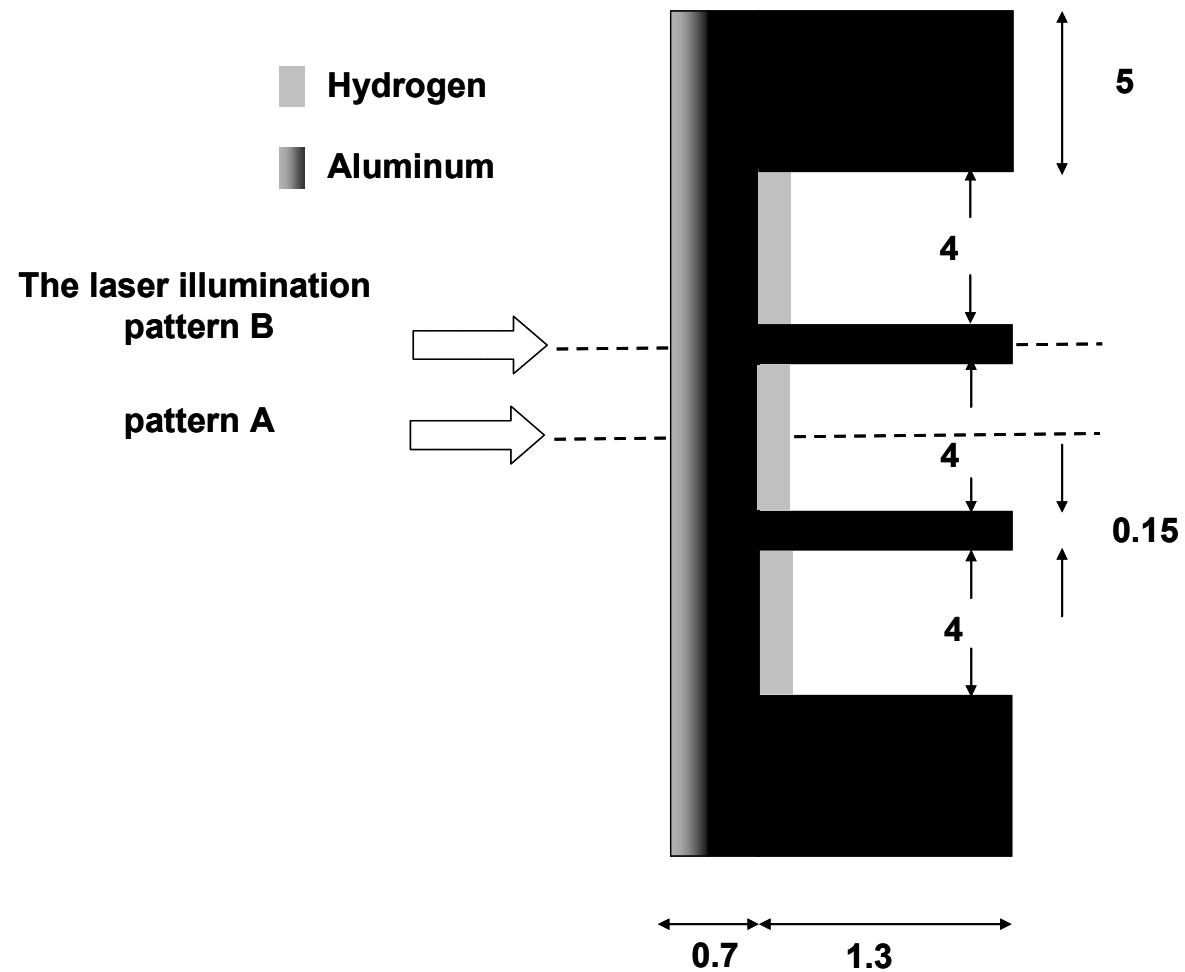


Proton Distribution in X-Y Space

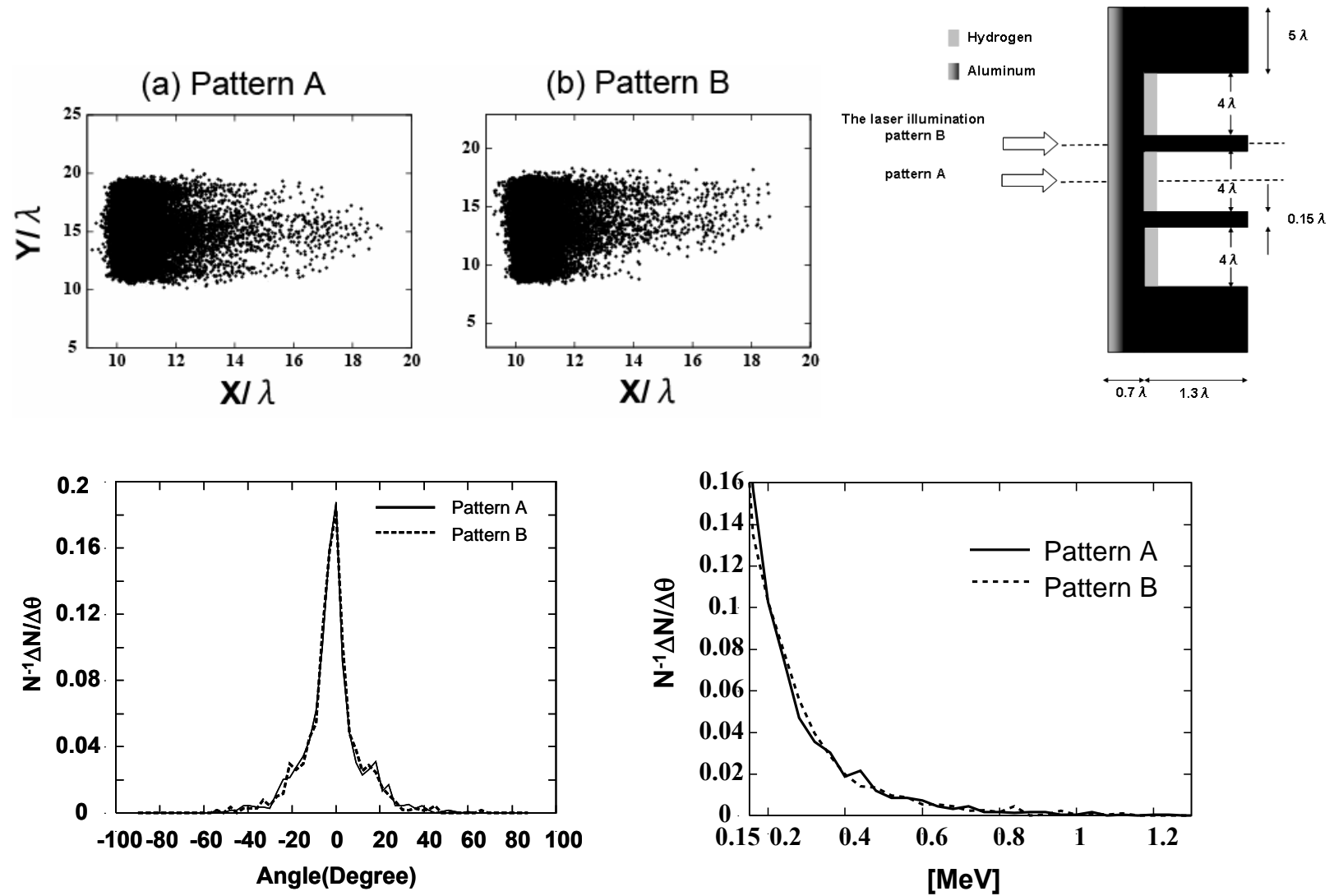


The proton beam transverse divergence of CaseB is suppressed successfully by the shaped target and the electron cloud localization.

Robust hole target against laser alignment



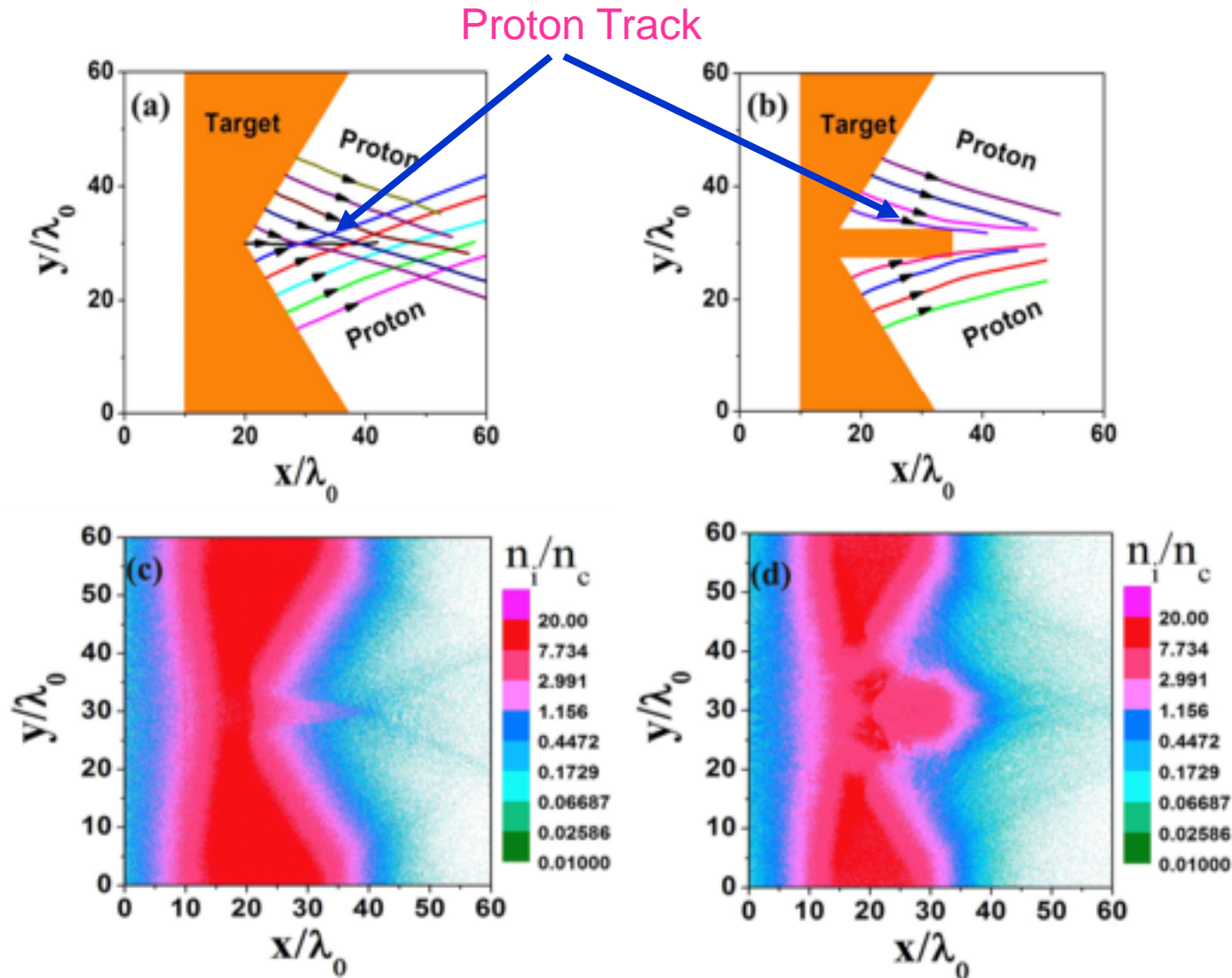
Robust hole target against laser alignment



Collimated proton beam from umbrella-like target

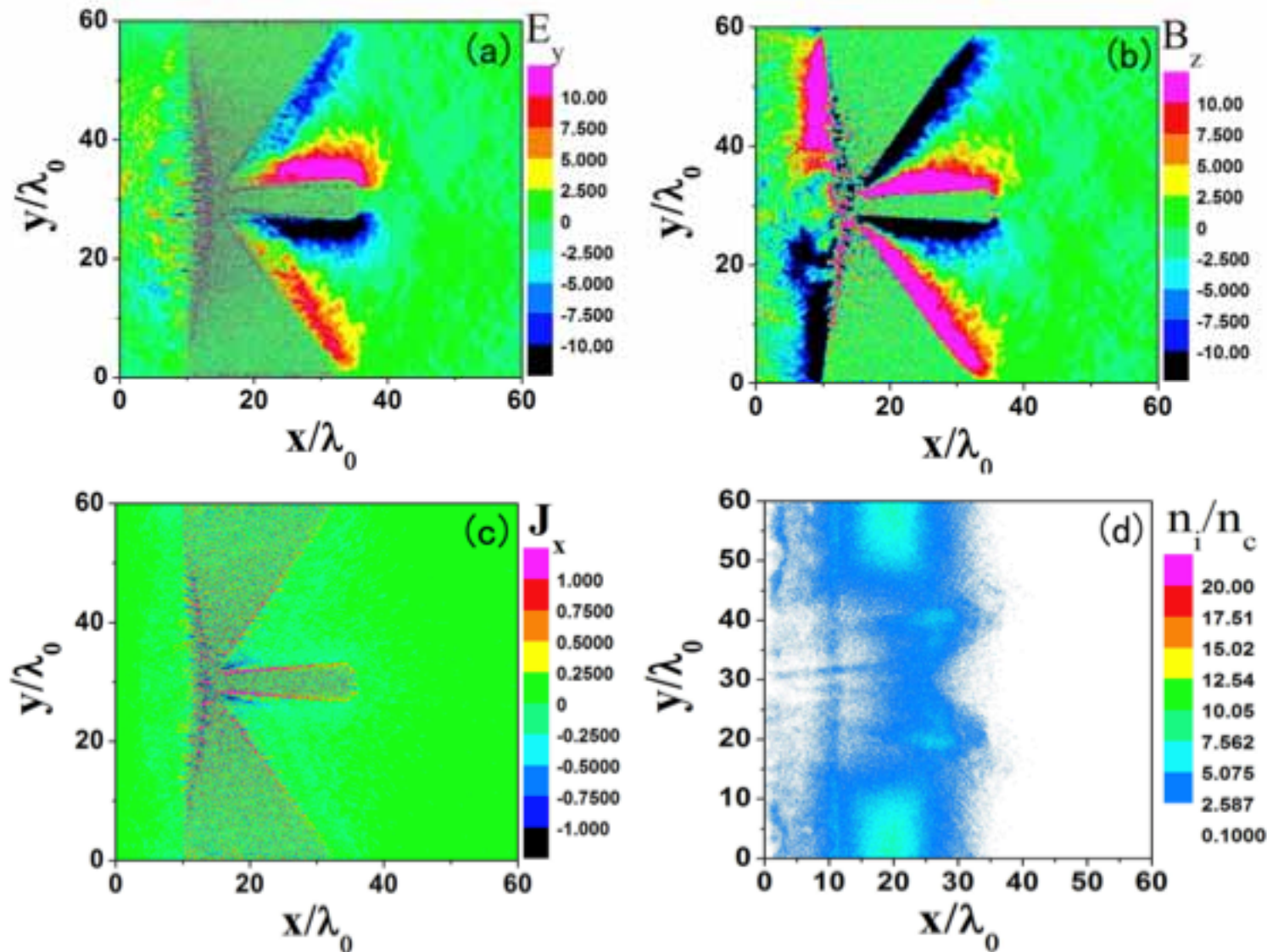
Y.Y. Ma, et al., Physics of Plasmas, 16(3), 034502, 2009.

Conic cavity target vs umbrella-like target



The wall : $8n_c$
filament : $12n_c$
 $I=10^{20}\text{W/cm}^2$
 $\tau=10\text{T}$

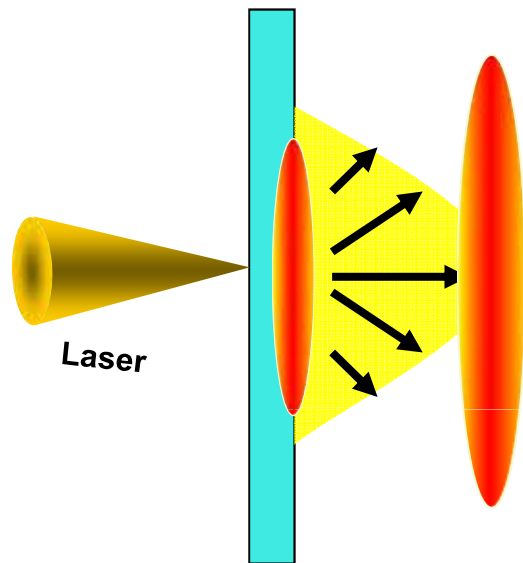
Umbrella-like target



208T

Energy Efficiency Enhancement

From Laser to Ion Beam

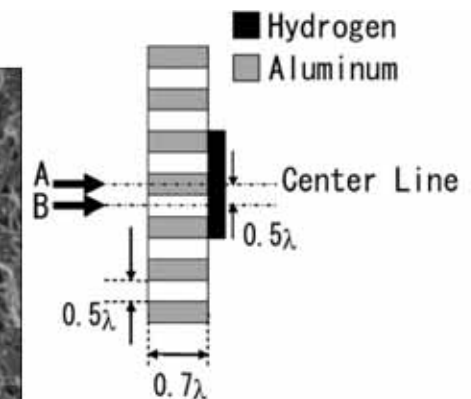
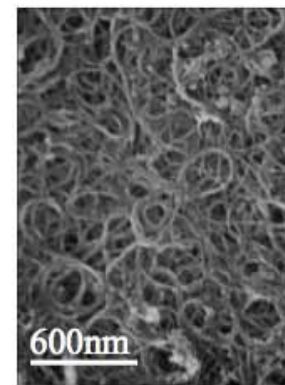
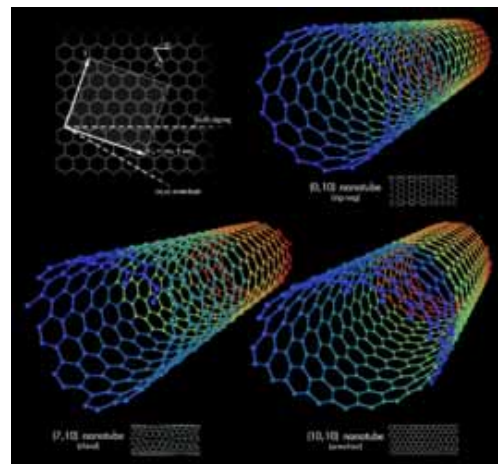
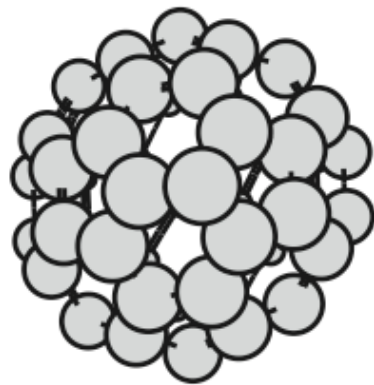


2. Low energy efficiency from laser to ion beam

~ a few % or less

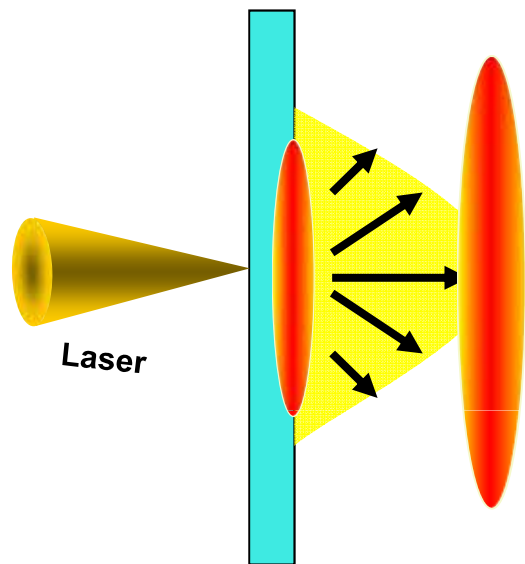


/ jaggy or rough surface target
/ cluster, amorphous, many-holes. ...
/ long life of electrons

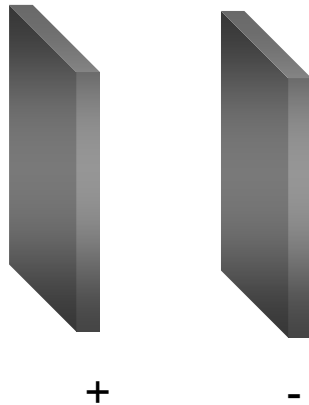


Energy Efficiency Enhancement

From Laser to Ion Beam



Ambipolar Field
~ Ion Diode



Virtual
Cathode



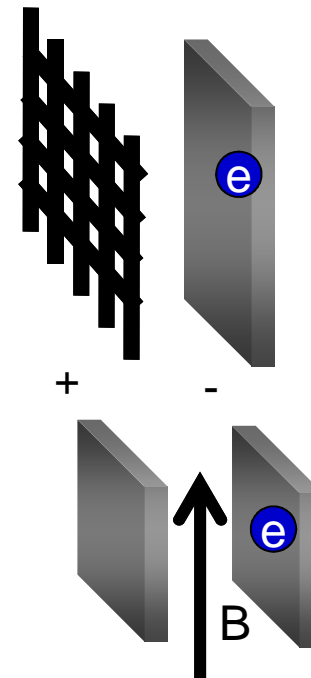
2. Low energy efficiency from laser to ion beam

~ a few % or less



/ long life of electrons

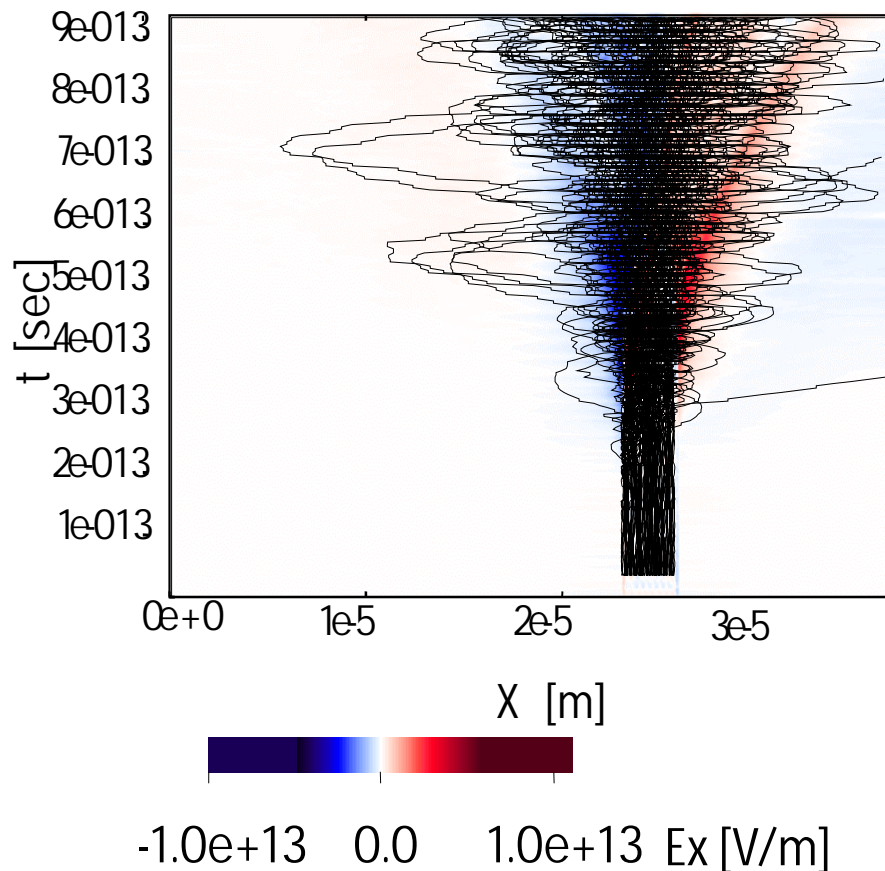
→ long life of acceleration E-field



Ion Acceleration

Target density $n_F = 2 \times n_c = 2.01 \times 10^{21} [cm^{-2}]$

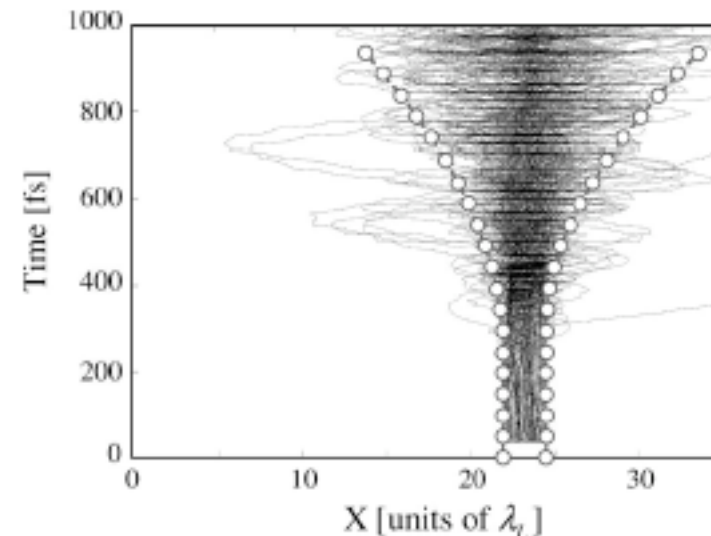
Electron trajectory



Electrons oscillate in a potential well of the ESF with a high frequency

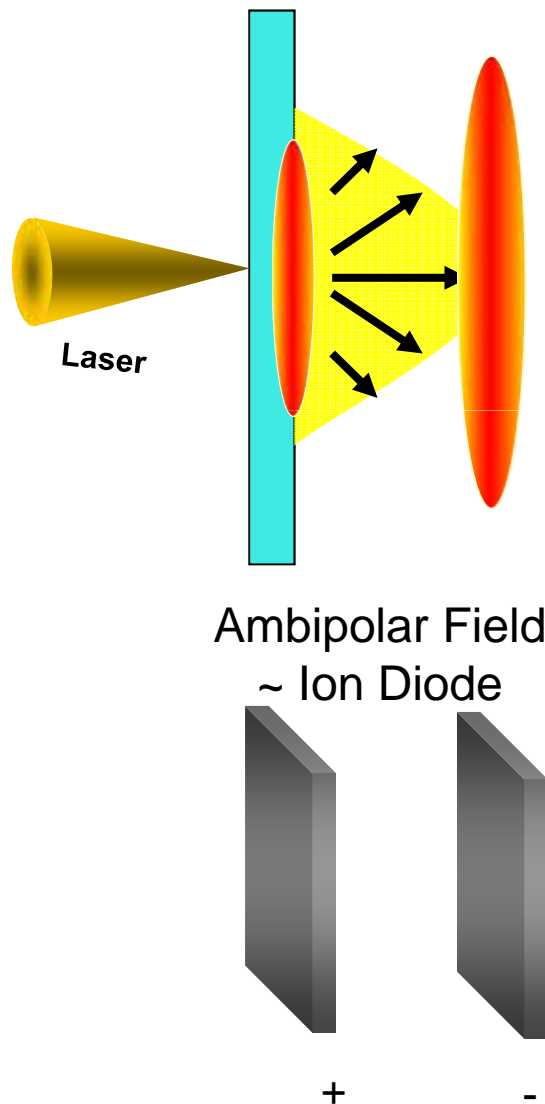


The number of the electrons to sustain the ESF wave depends on the energy of the electrons



Energy Efficiency Enhancement

From Laser to Ion Beam

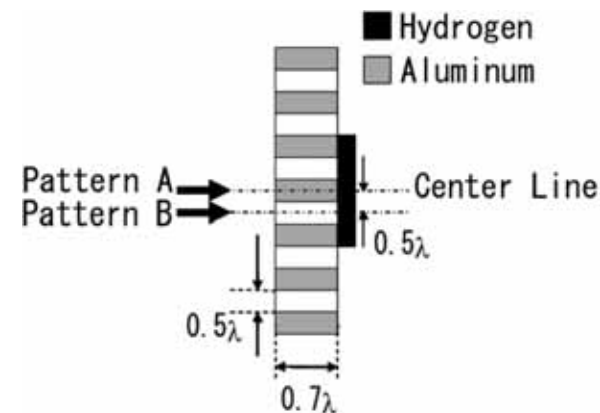
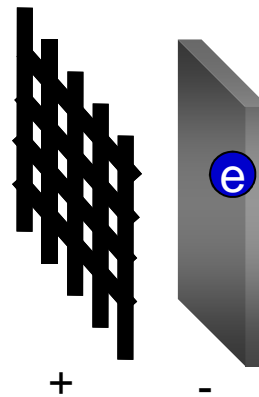


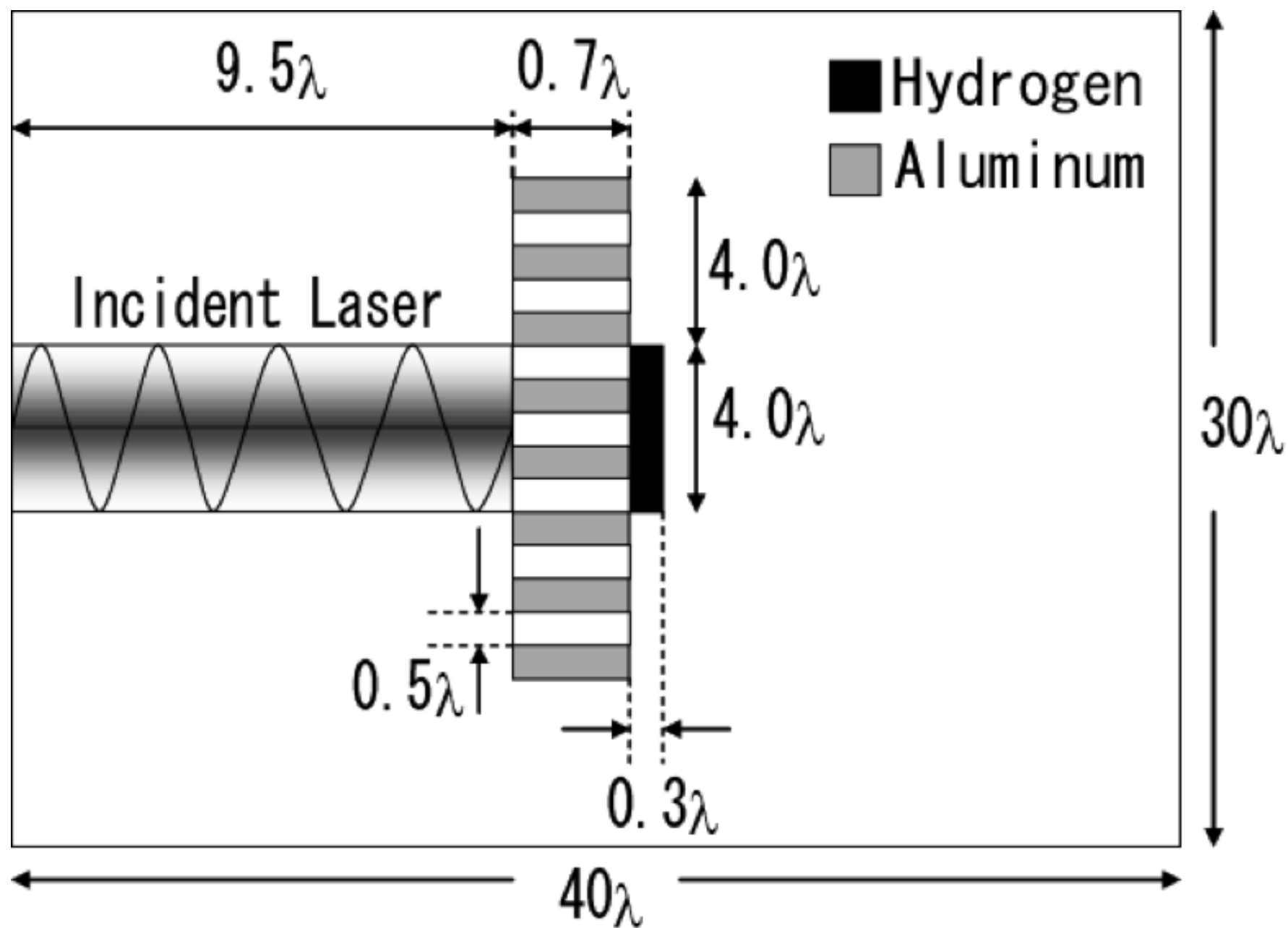
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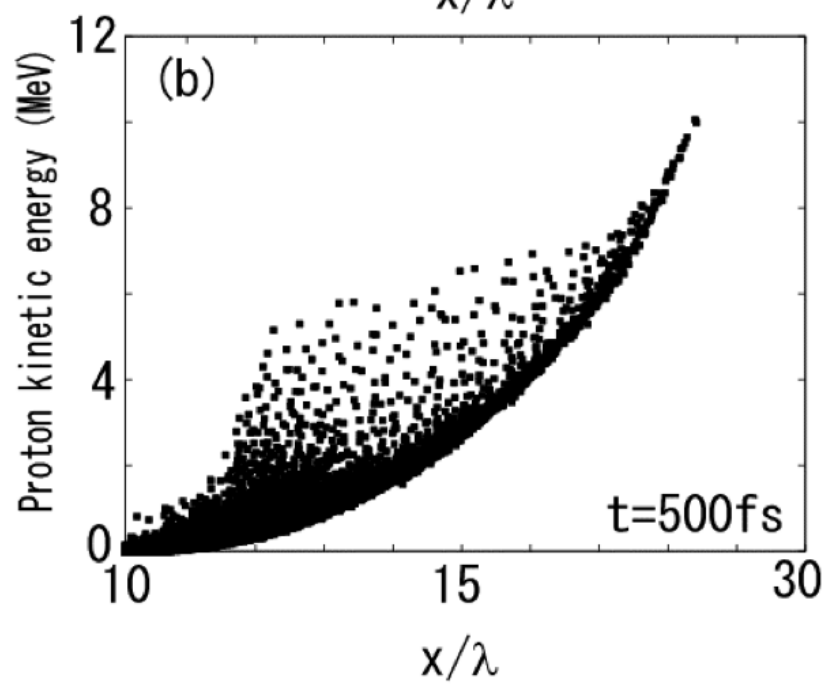
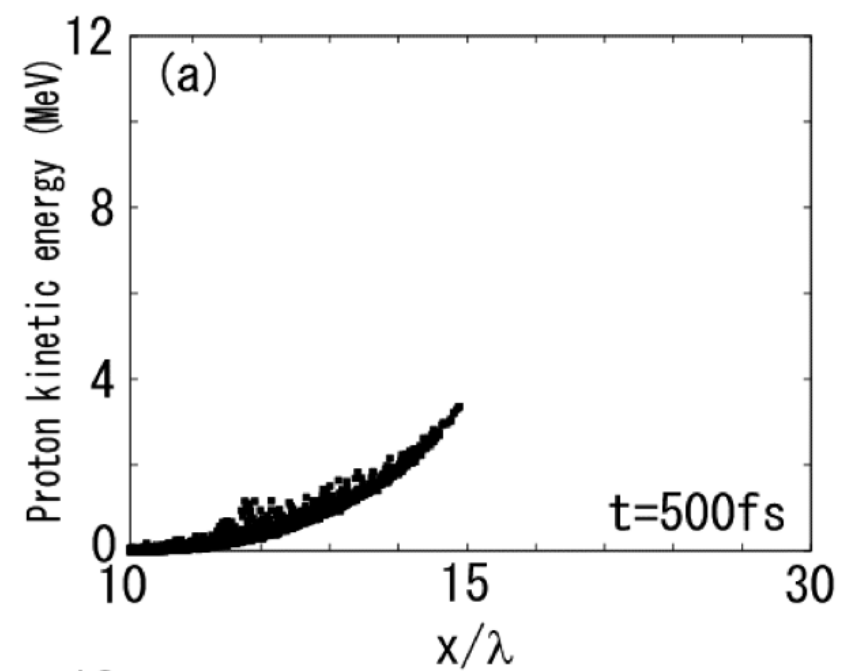
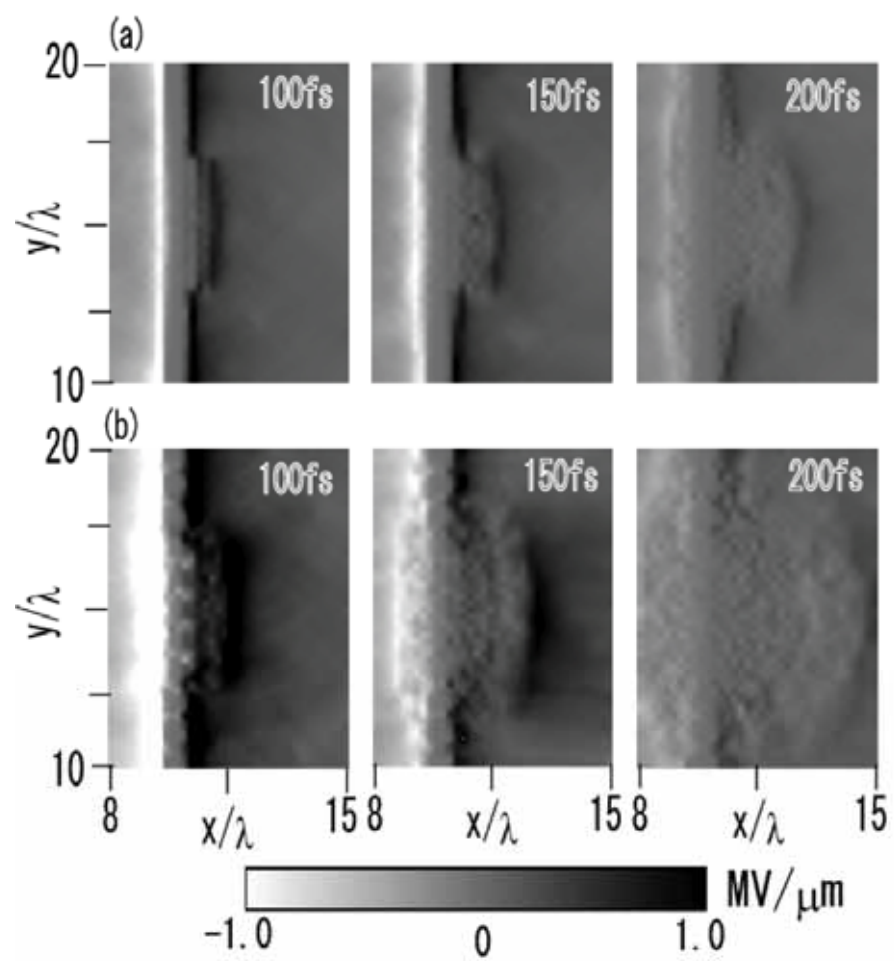
~ a few % or less

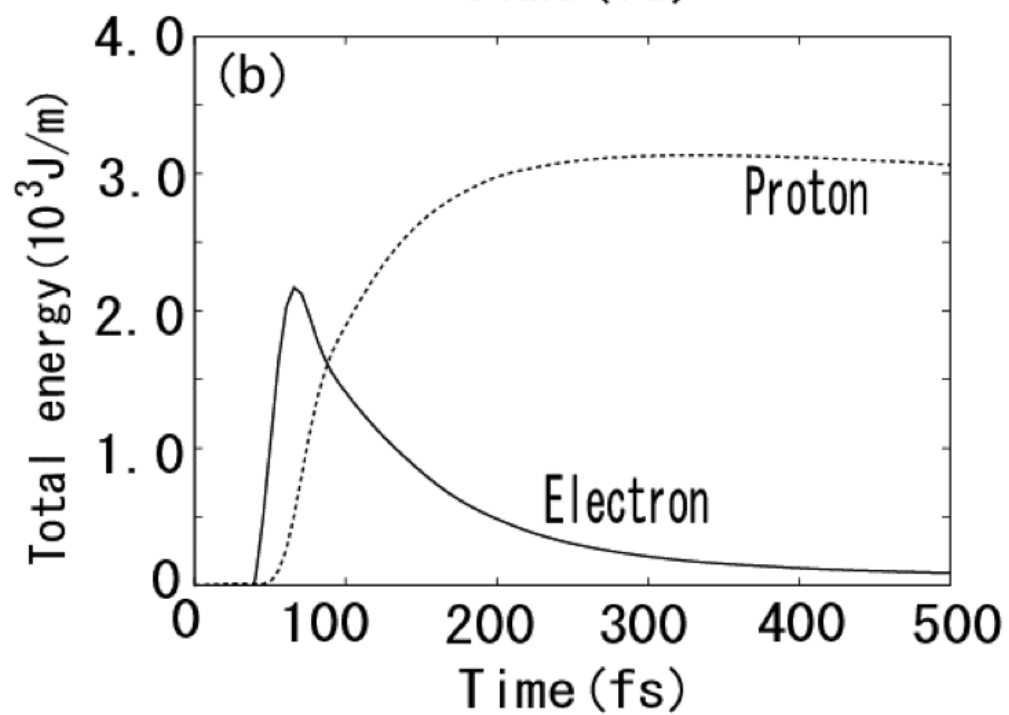
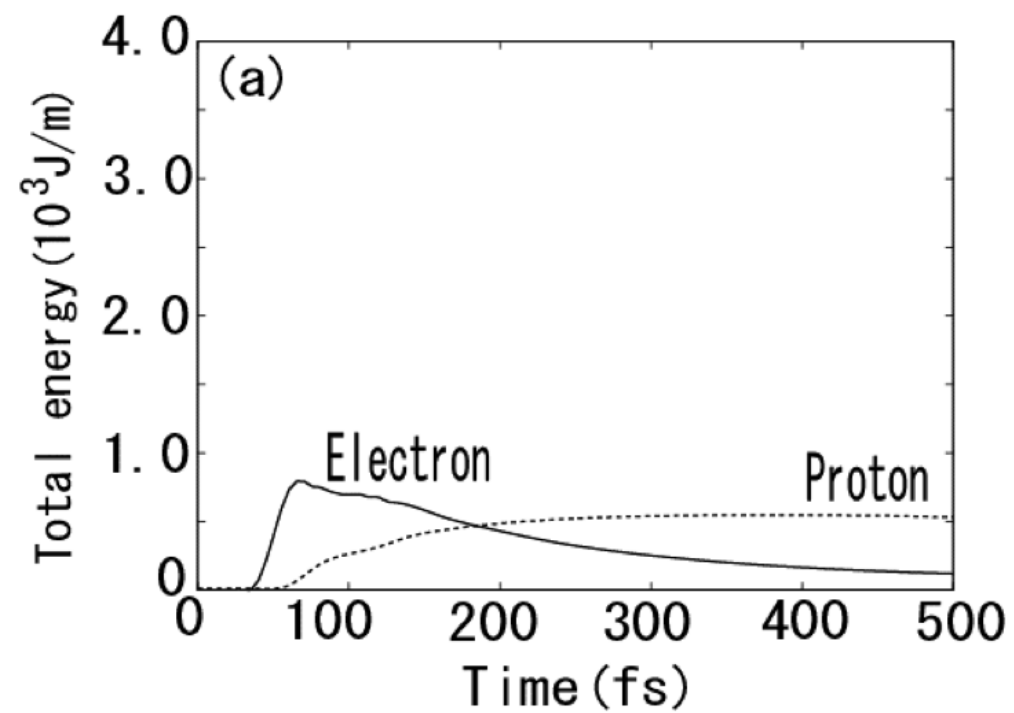


/ long life of electrons









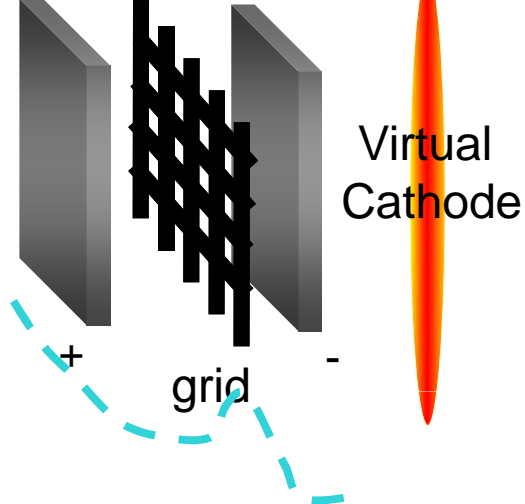
Problems of Laser Ion Accelerator

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Energy Spectrum Control

Ambipolar Field

~ Ion Diode



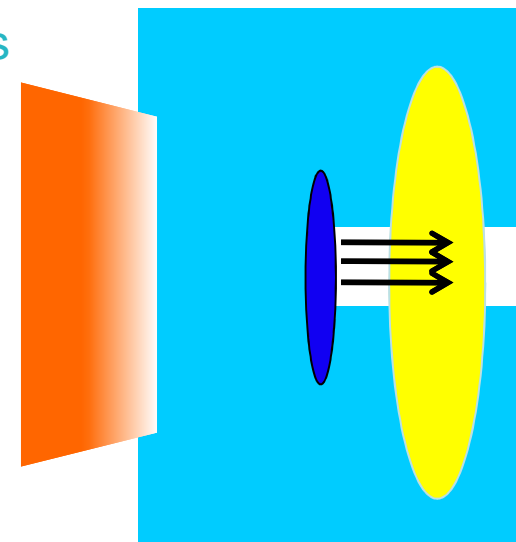
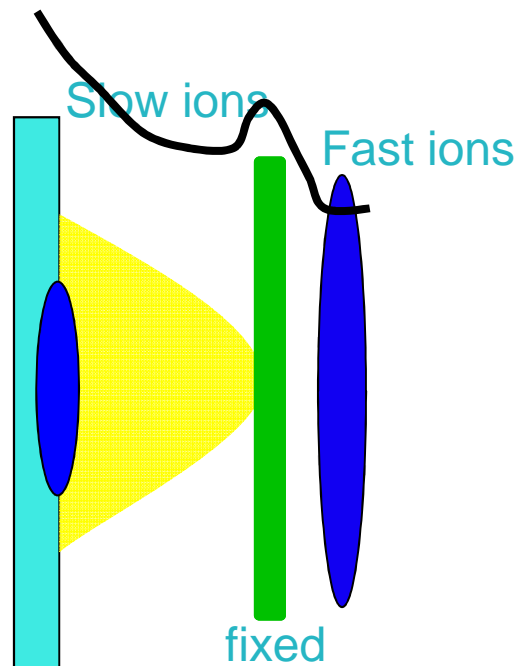
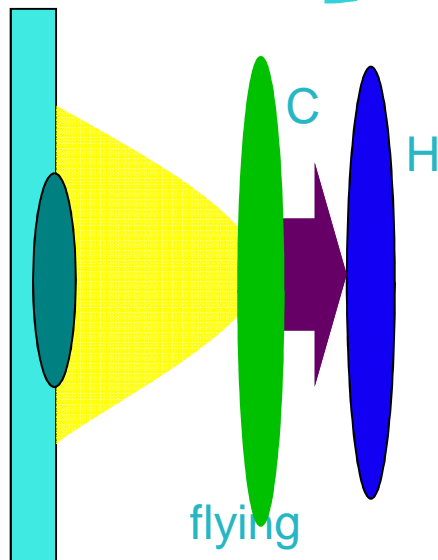
1. Ion beam quality

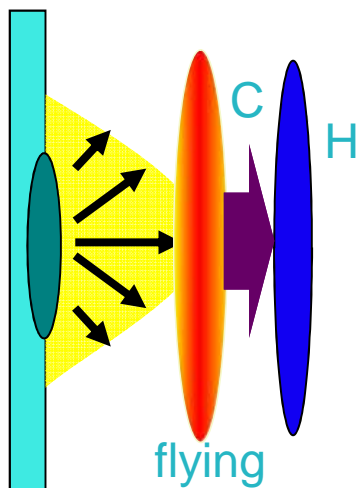
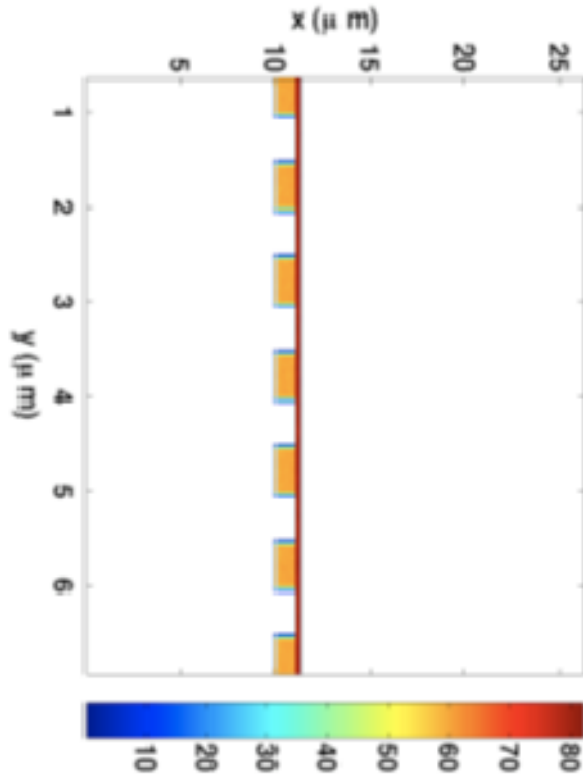
--- energy spectrum control



/ very thin ion source < skin depth

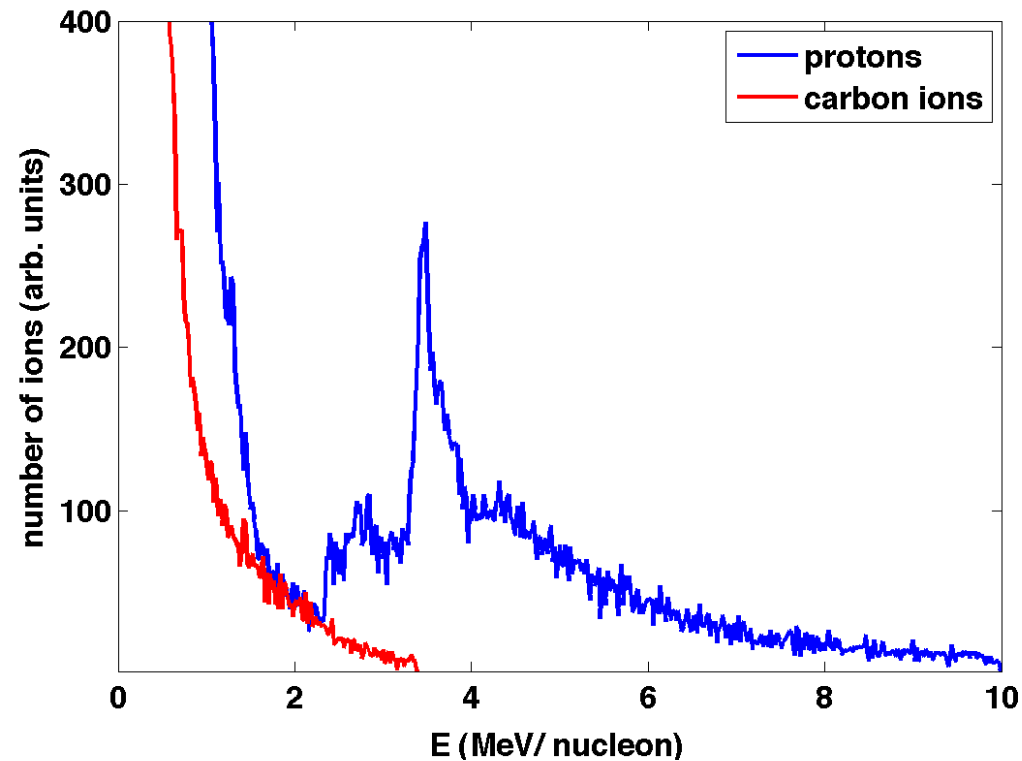
/ double layer – flying or fixed





Laser: intensity 10^{20} W/cm²
 wavelength 1 μ m
 duration 15 fs
 pulse shape \sin^2
 polarization P
 spot size 3 μ m
 incidence angle 0

Foil: thickness 300 nm
 density 0.4 g/cm³
 composition **CH2**
 ion charge state C6+, H+ initial
 temperature 1 eV
 profile step-like boundary conditions
 thermalization of fast electrons



case 3 periodic array with structure composed of boxes of 60 nm × μm separated by 0.25 μm, material composition C⁶⁺, density 0.35 g/cm³

= > Flat target: 2.8% (Laser -> ions)
Structured target: 26%(Laser -> ions) case 3(20% C, 6%Protons)

Quasi-monoenergetic proton beam from target with holed backside

T. P. Yu, Y.Y. Ma, et al., Physics of Plasmas, 16(3), 033112, 2009.

The production of quasi-monoenergetic protons from microstructured targets

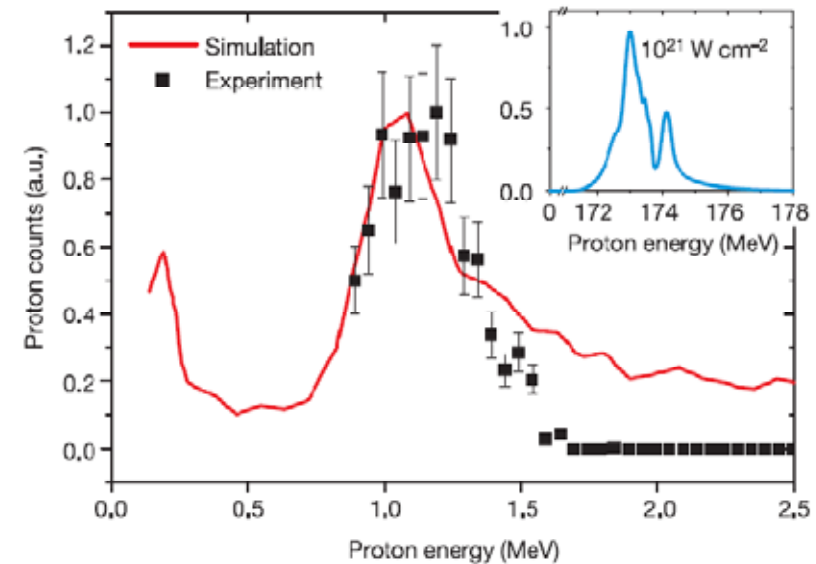
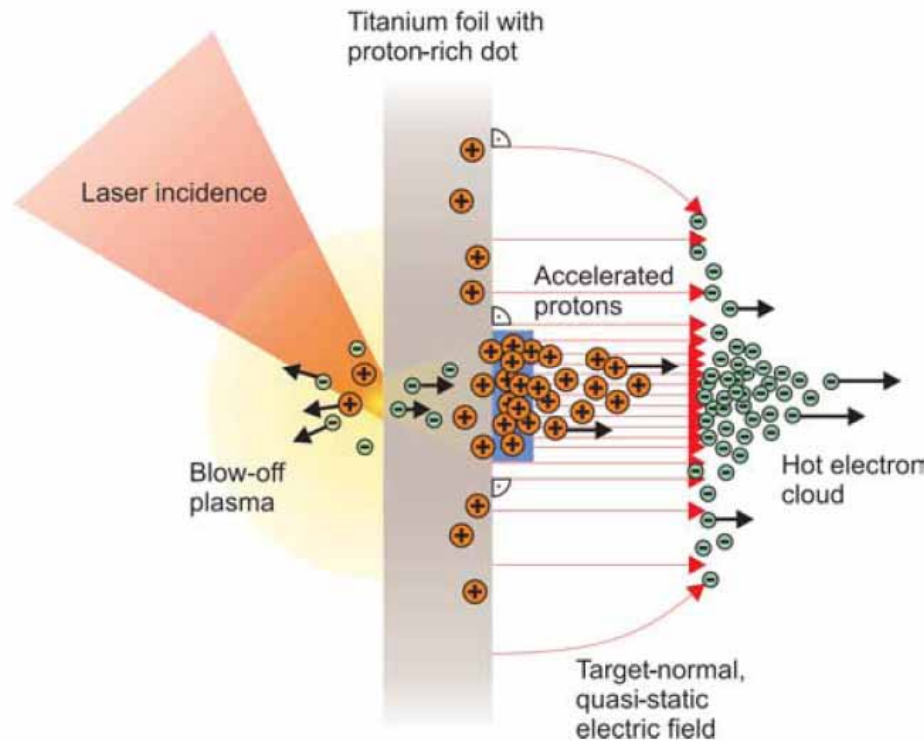


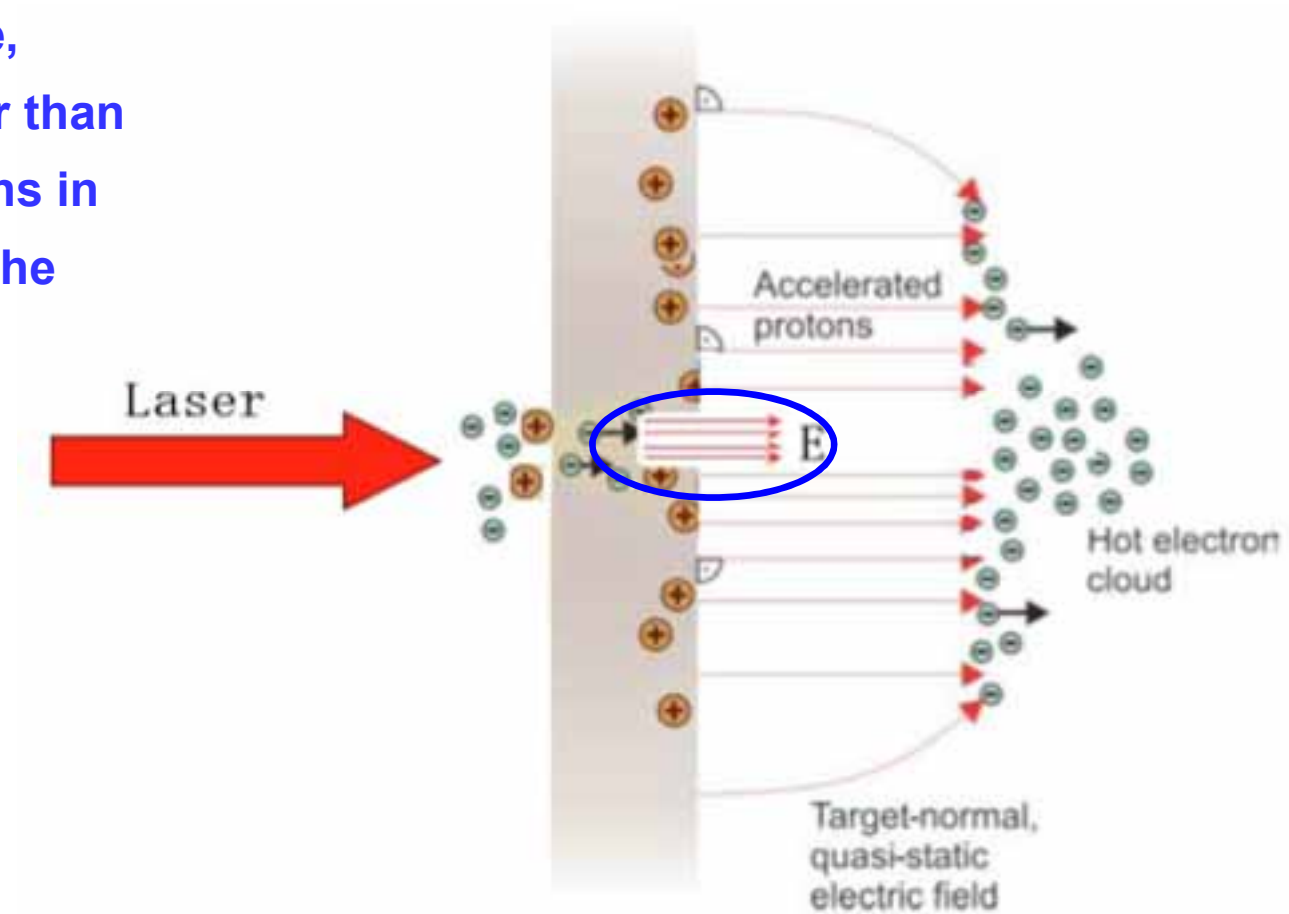
Figure 4 | Results from simulations and scalability of the technique. Comparison of experimental data (black squares) to the proton spectrum obtained from two-dimensional-PIC simulation (red line) for following conditions: laser intensity $I_L = 3 \times 10^{19} \text{ W cm}^{-2}$, and target dimensions $5 \mu\text{m Ti foil} + 0.5 \mu\text{m PMMA dot } (20 \times 20) \mu\text{m}^2$. Experimental data points comprise the observable energy range on the MCP detector. The statistical uncertainty for the measured data has a value of 20% s.d. as shown by the error bars. In the inset a simulation for a petawatt-laser system demonstrating the scalability of proton acceleration from microstructured targets is shown. The parameters for the simulation are $I_L = 1.2 \times 10^{21} \text{ W cm}^{-2}$, $5 \mu\text{m Ti foil} + 0.1 \mu\text{m PMMA dot } (2.5 \mu\text{m diameter})$. The proton spectrum exhibits a narrow peak with relative energy width of $\Delta E/E \approx 1\%$ at a peak energy of 173 MeV.

H. Schworer, Nature 439, 445(2006)

German

The Physical model for micro-hole target

A micro-hole is made, whose size is smaller than the laser spot. Protons in the hole almost feel the unique electric field.



PIC simulation parameters

FIG 1. Scheme of the simulation geometry. The target is located from $x=5.0\lambda_0$ to $x=7.0\lambda_0$.

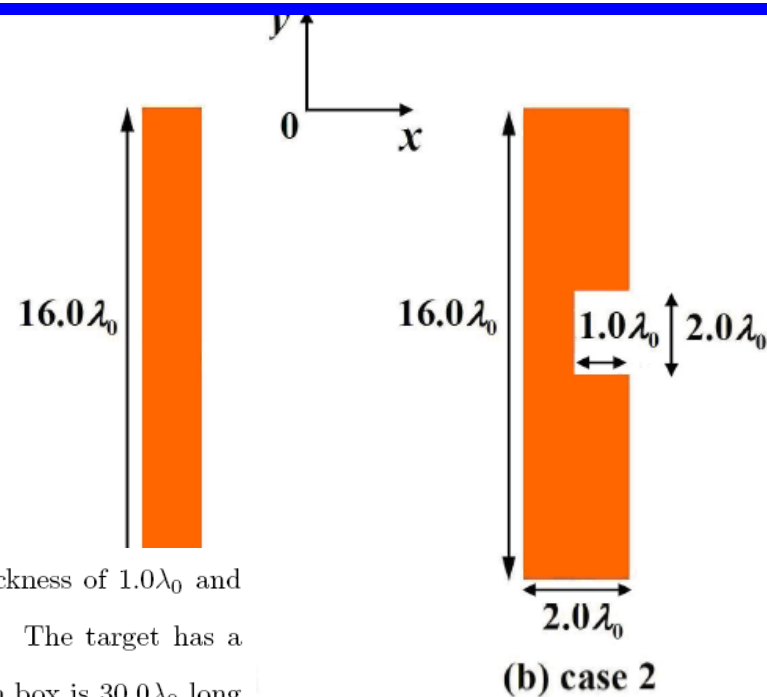


FIG. 1: (Color online.) (a) Case 1: simple flat target. The target has a thickness of $1.0\lambda_0$ and is located between $x = 5.0\lambda_0$ and $6.0\lambda_0$. (b) Case 2: holed-backside target. The target has a thickness of $2.0\lambda_0$ and is located between $x = 5.0\lambda_0$ and $7.0\lambda_0$. The simulation box is $30.0\lambda_0$ long and $16.0\lambda_0$ wide. The laser impinges on the target from the left.

Laser parameters : $I = 1.0 \times 10^{21} \text{ W / cm}^2$ $\lambda_0 = 1.06 \mu\text{m}$ $R = 10 \mu\text{m}$

$$\tau = 10.0 T_0 (\approx 33.3 \text{ fs}) \quad a = a_0 \sin^2(\pi t / \tau_L) \exp(-y^2 / 2R^2)$$

Plasm parameters : $n_0 = 20n_c$ $t_e = 1000 \text{ eV}, t_i = 333 \text{ eV}$ $h = 2.0\lambda_0, d = 1.0\lambda_0$

T. P. Yu, Y.Y. Ma, et al., Physics of Plasmas, 16(3), 033112, 2009.

compare

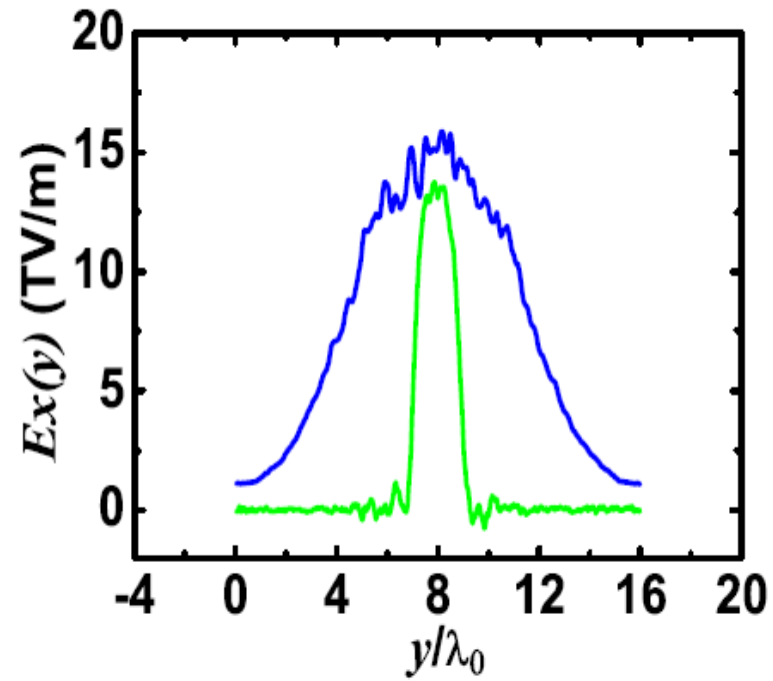


FIG. 2: (Color online.) Transverse distribution of the longitudinal sheath electric field E_x at $t = 16.08T$. Laser focal radius is $R = 2.0\lambda_0$ for Case 1 (blue/dark) and Case 2 (green/bright).

Energy spectra

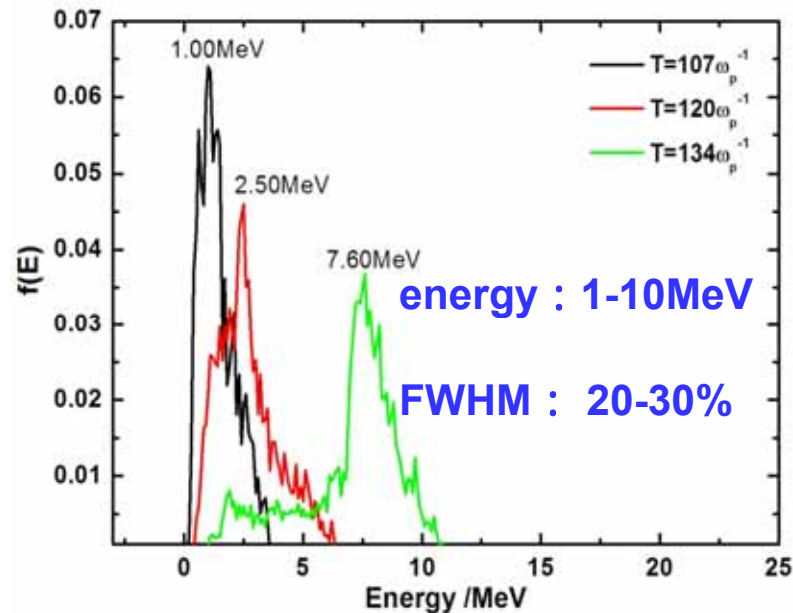


FIG 7. Proton energy spectra obtained from the concave at $t=107\omega_p^{-1}$ (black), $t=120\omega_p^{-1}$ (red) and $t=134\omega_p^{-1}$ (green)

Future application to Ion Therapy

1. Doctor & Patient – friendly system for individual patient
2. Optimization of Curing Plan – Which direction, how much particle energy, how many times,
3. Ion illumination & deposition code OK2 + parallelization
4. Integrated PSE System
5. Compact ion accelerator – Laser ion accelerator

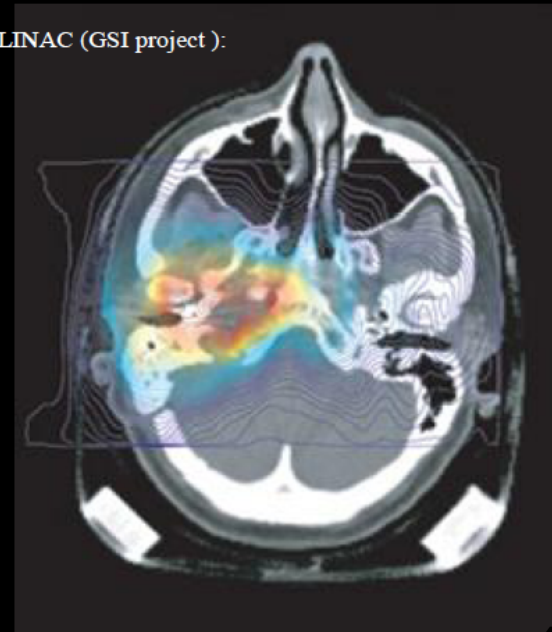
Code OK2 – one possibility of medical application

Output beam parameters from carbon LINAC (GSI project):

Design ion	$^{12}\text{C}^{4+}$
Beam energy	7 MeV/u
Beam pulse	200 μs
Beam emittance	1,88 mm-mrad

OK2 can be combined with PET method as an additional tool:

- 1) For adjusting the first beam shot precisely.
- 2) For dose calculation accumulated in the tumor area as well as in the surrounding tissues in time of therapeutic irradiation.
- 3) For optimization of a patient position before irradiation.
- 4) For simulation of multi-beam irradiation on the tumor area.



Thank you!

Location



From **Narita**
(Tokyo)
International
Airport

-> 3hours by
Bus or Train

From **Tokyo**
1hour by Train
(Shinkan-sen)

Nikko – 日光 – one of World Heritages

